

18th Synthesis Imaging Workshop

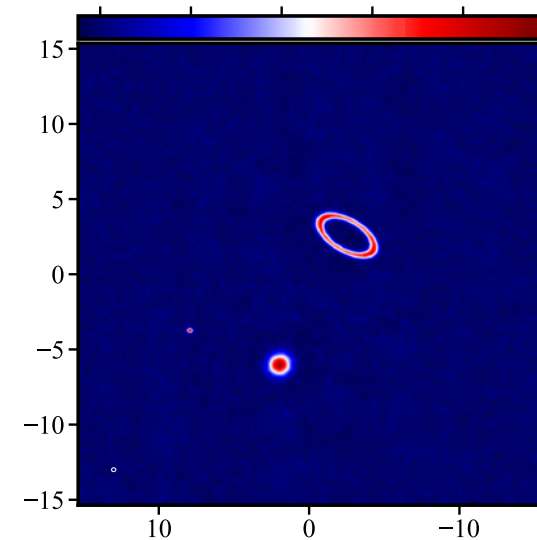
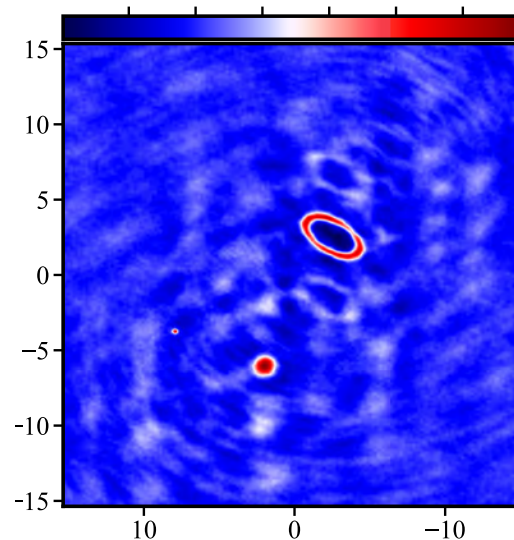
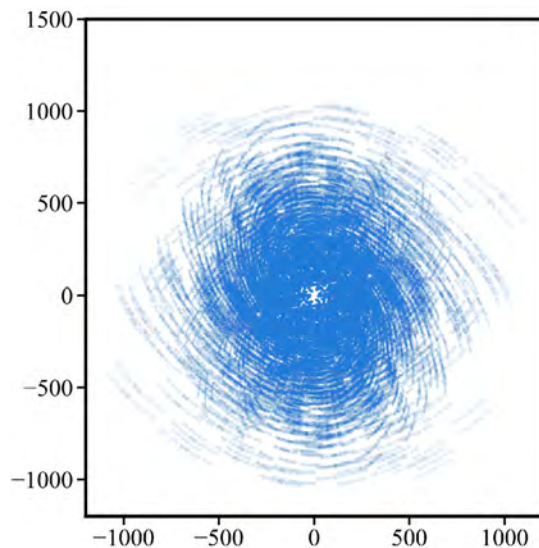
May 18, 2022—May 25, 2022

🕒 Viewing in Mountain Time [Adjust](#)
Virtual



Imaging Basics

David J. Wilner CENTER FOR **ASTROPHYSICS**
HARVARD & SMITHSONIAN



Overview

goals

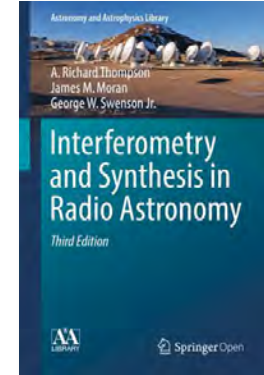
- gain intuition about interferometric imaging
- understand the need for deconvolution

topics

- get comfortable with Fourier Transforms
- review visibility concept and (u,v) plane sampling
- formal description of imaging
- imaging in practice: FFT, gridding, weighting schemes
- deconvolution and the clean algorithm

References

- Thompson, A.R., Moran, J.M. & Swensen, G.W. “Interferometry and Synthesis Imaging in Radio Astronomy” 3rd edition, 2017, Springer Open
- NRAO Synthesis Workshop proceedings
 - Synthesis Imaging in Radio Astronomy II,, ASP Conference Series, Vol. 180, 1999, eds. Taylor, G.B., Carilli, C.L., Perley, R.A.
 - lecture slides: www.aoc.nrao.edu/events/synthesis
- IRAM 2000 Interferometry School proceedings
 - www.iram.fr/IRAMFR/IS/IS2008/archive.html
- Condon, J.J. & Ransom, S.M. 2016 “Essential Radio Astronomy”
 - science.nrao.edu/opportunities/courses/era
- Many useful pedagogical presentations available on-line
 - ALMA videos, casadocs, CASS radio school, ERIS lectures, ...
- EHT M87 imaging paper, 2019, ApJ Letters, 875, L4



Visibility and Sky Brightness

For small field of view, far field, quasi-monochromatic, incoherent, etc.:

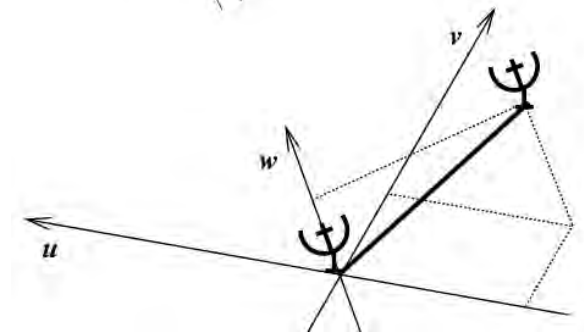
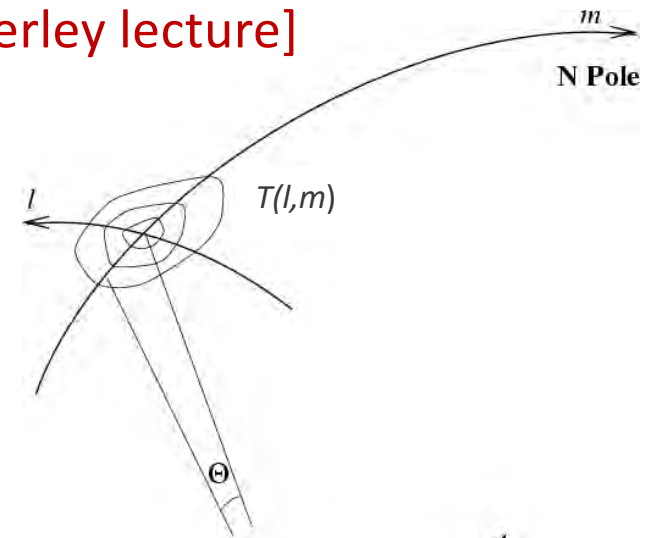
*The complex visibility function $V(u,v)$ is the **2D Fourier Transform** of the sky brightness distribution $T(l,m)$* [Rick Perley lecture]

$$V(u, v) = \int \int T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

(u, v) are (E-W, N-S) spatial frequencies [wavelengths]

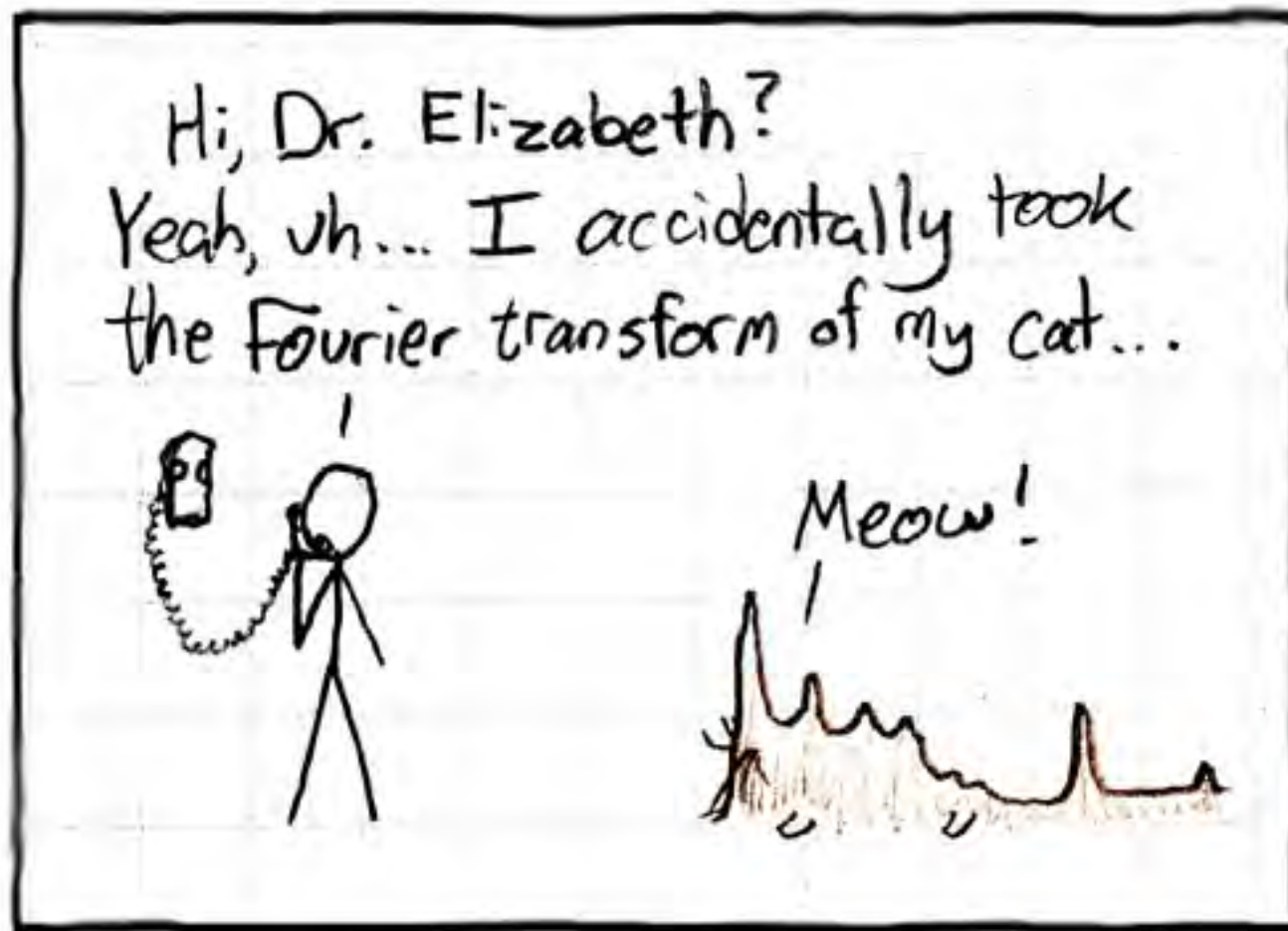
(l, m) are (E-W, N-S) angles in the tangent plane [radians]

recall $e^{ix} = \cos x + i \sin x$



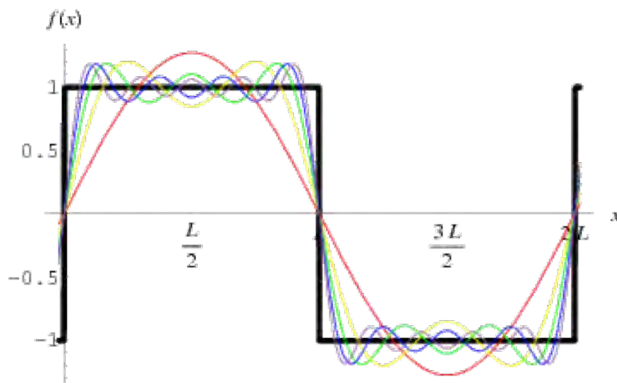
**Jean Baptiste
Joseph Fourier
1768-1830**

xkcd.com/26/



Fourier Fundamentals

- acquire comfort with the Fourier domain...
 - functions and their Fourier transforms occupy *upper* and *lower* domains, as if “functions circulated at ground level and their transforms in the underworld” (Bracewell 1965)
- Fourier theory states that any well behaved signal can be represented as the sum of sinusoids



$$f(x) = \frac{4}{\pi} \sum_{n=1,3,5,\dots}^{\infty} \frac{1}{n} \sin\left(\frac{n\pi x}{L}\right).$$

Wolfram Mathworld



- the Fourier transform is the mathematical tool that decomposes a signal into its sinusoidal components (frequency, amplitude, phase)
- the Fourier transform contains *all* information of the original signal

Useful Fourier Transform Properties

- **addition** in one domain is **addition** in the other

$$g(x) + h(x) = G(s) + H(s)$$

- **multiplication** in one domain is **convolution** in the other

$$g(x) = h(x) * k(x) \quad G(s) = H(s)K(s)$$

- “**large**” in one domain is “**small**” in the other

$$g(\alpha x) = \alpha^{-1} G(s/\alpha)$$

- an **offset** in one domain is a **phase shift** in the other

$$g(x - x_0) = G(s) e^{i2\pi x_0 s}$$

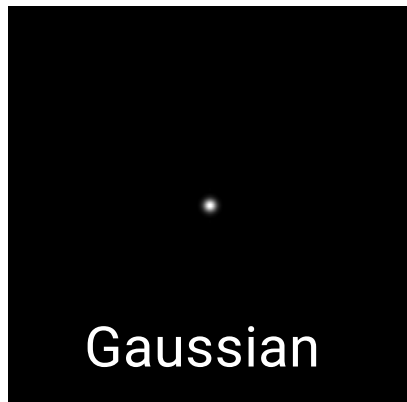
- the Nyquist-Shannon **sampling theorem**

$$g(x) \subset \Theta \text{ completely determined if } G(s) \text{ sampled at } \leq 1/\Theta$$

Some 2D Fourier Transform Pairs

$T(l,m)$

$V(u,v)$ amplitude



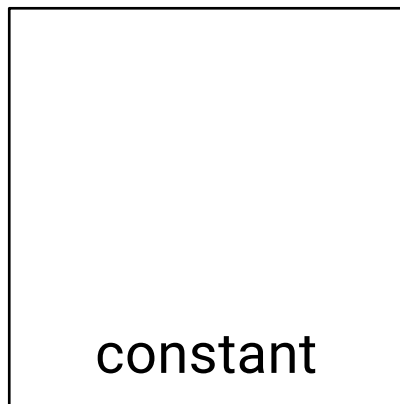
$\mathcal{F} \rightarrow$



*narrow features
transform into
wide features
(and vice-versa)*



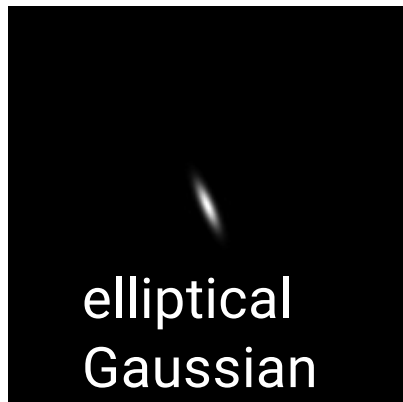
$\mathcal{F} \rightarrow$



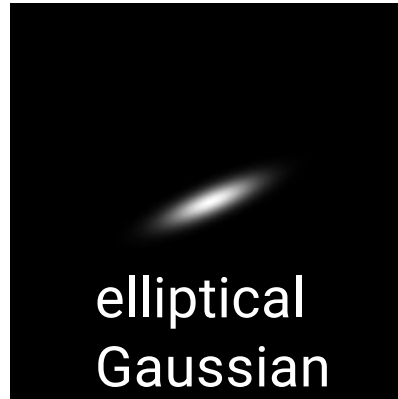
Some 2D Fourier Transform Pairs

$T(l,m)$

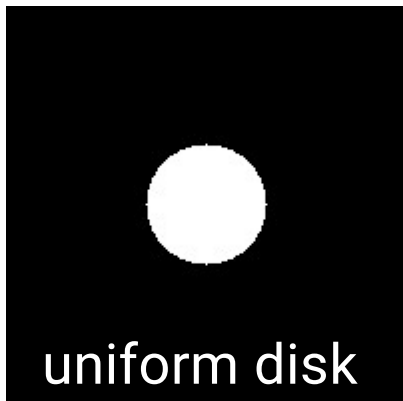
$V(u,v)$ amplitude



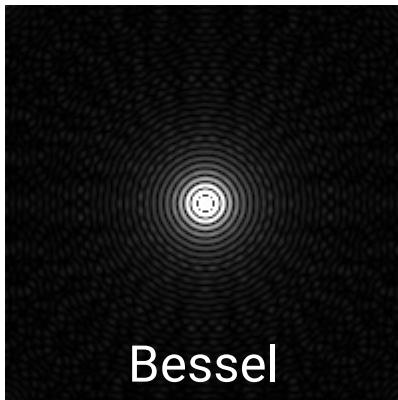
$\mathcal{F}$
→



*narrow features
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(and vice-versa)*



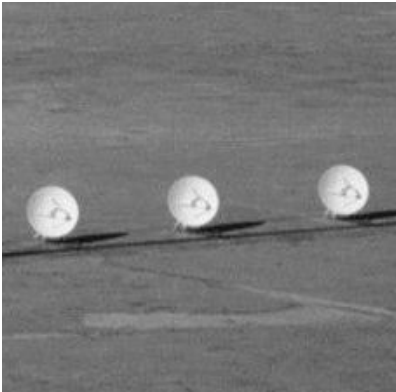
$\mathcal{F}$
→



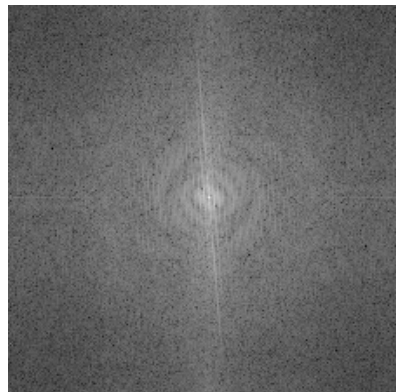
*sharp edges
result in many
high spatial
frequencies*

Amplitude and Phase

$T(l,m)$



$V(u,v)$ amplitude

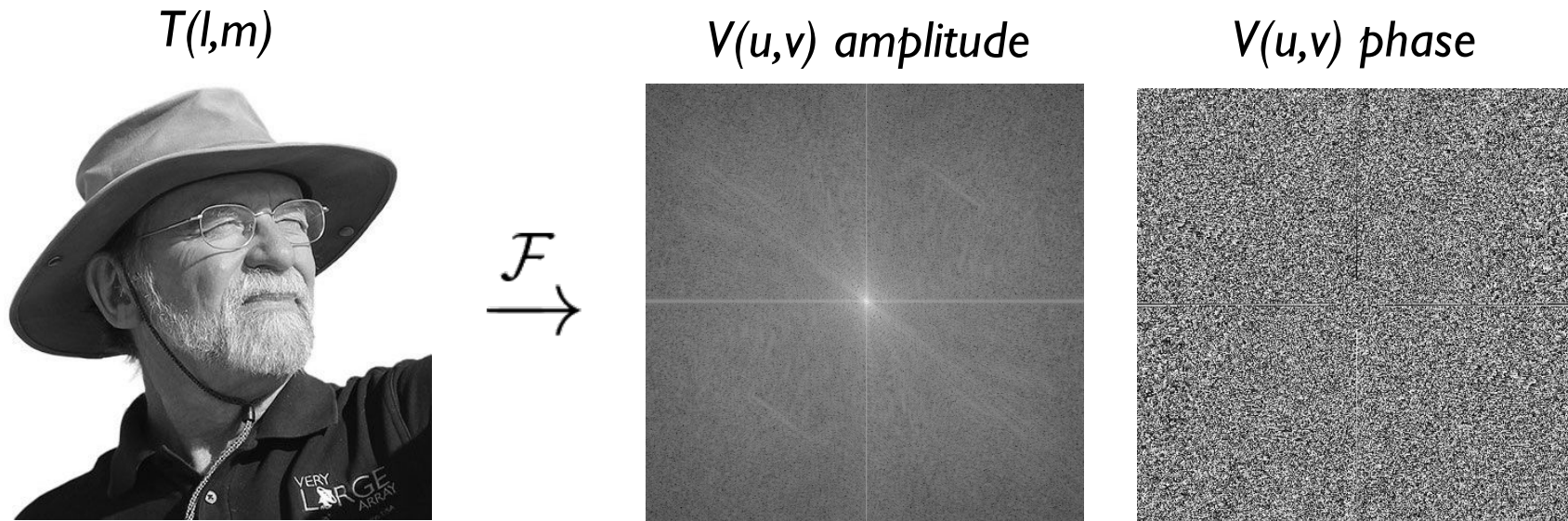


$\mathcal{F}$
→

amplitude tells “how much” of a certain spatial frequency

phase tells “where” the spatial frequency is located

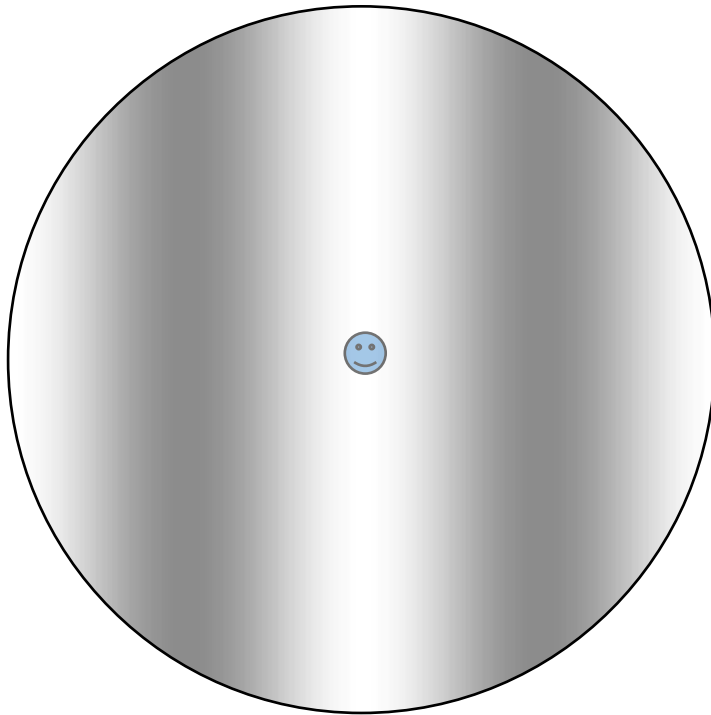
A more complicated example



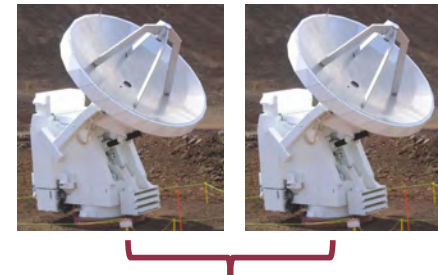
$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

- **$T(l,m)$ is real:** $V(-u,-v) = V^*(u,v)$ where $*$ = complex conjugate
 - get two visibilities for one measurement
- **$V(u=0,v=0)$ is the integral of $T(l,m)dl dm$ = total flux density**

$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

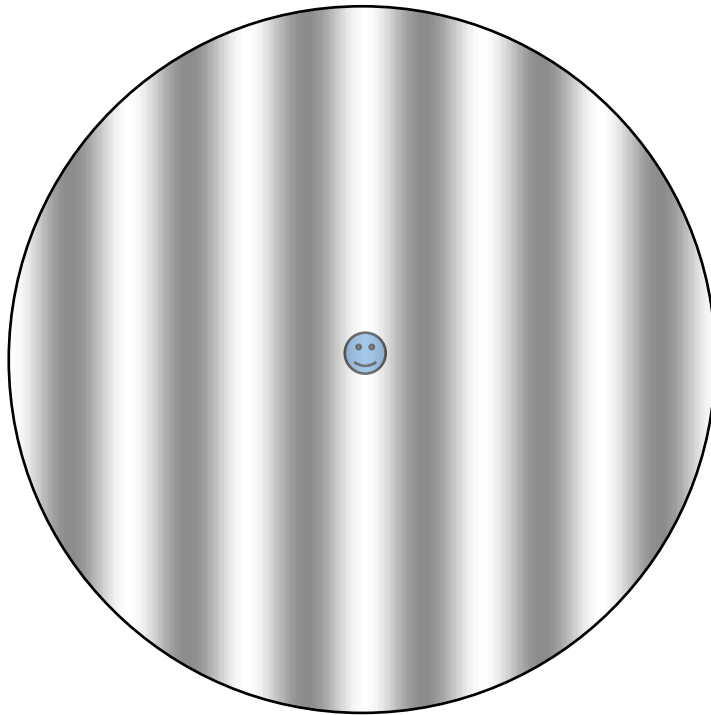


compact source



short baseline
wide fringe pattern

$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

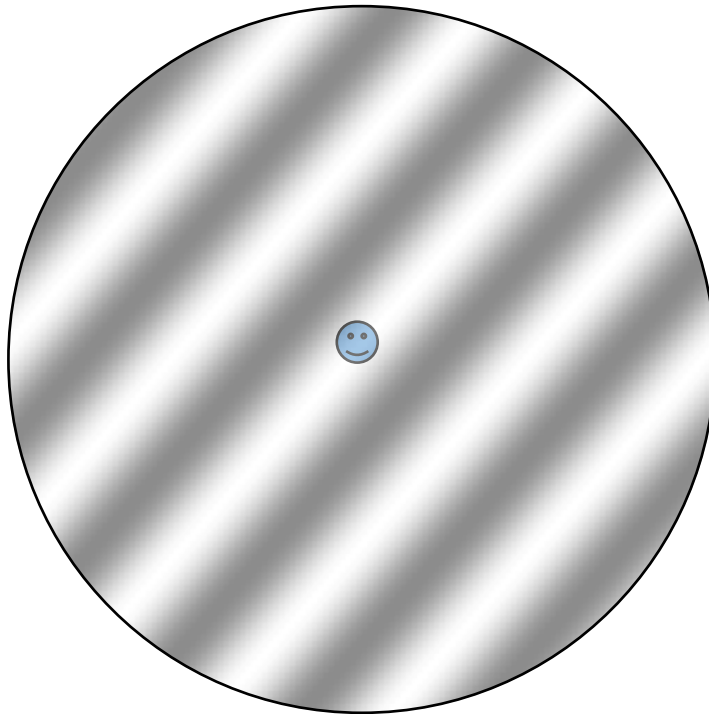


compact source

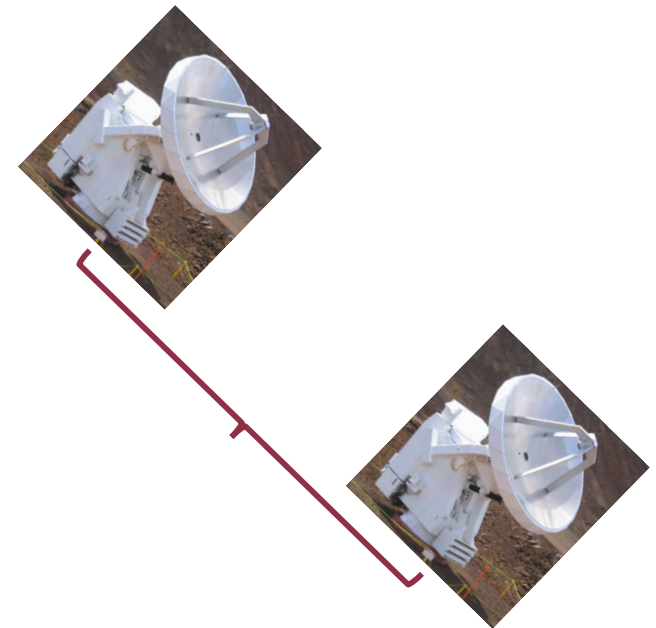


long baseline
narrow fringe pattern

$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

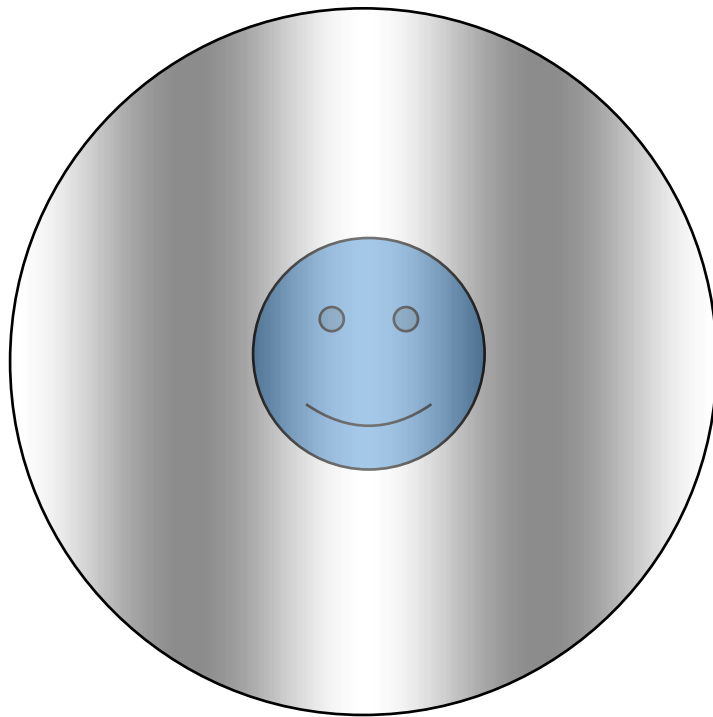


compact source

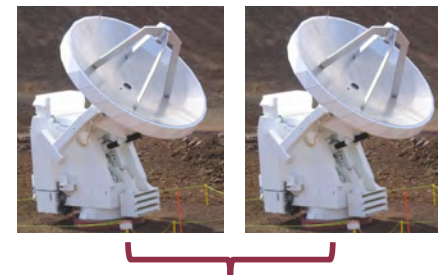


long baseline
narrow fringe pattern
different orientation

$$V(u, v) = \int \int T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

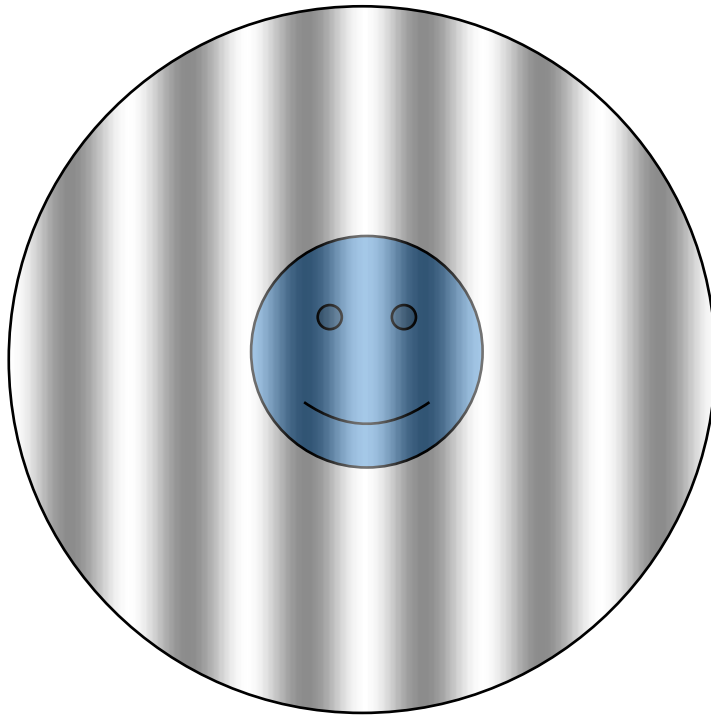


extended source

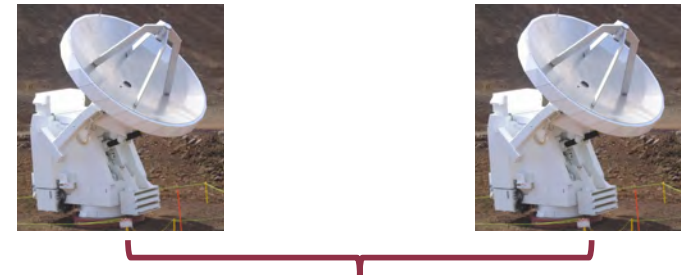


short baseline
wide fringe pattern

$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$

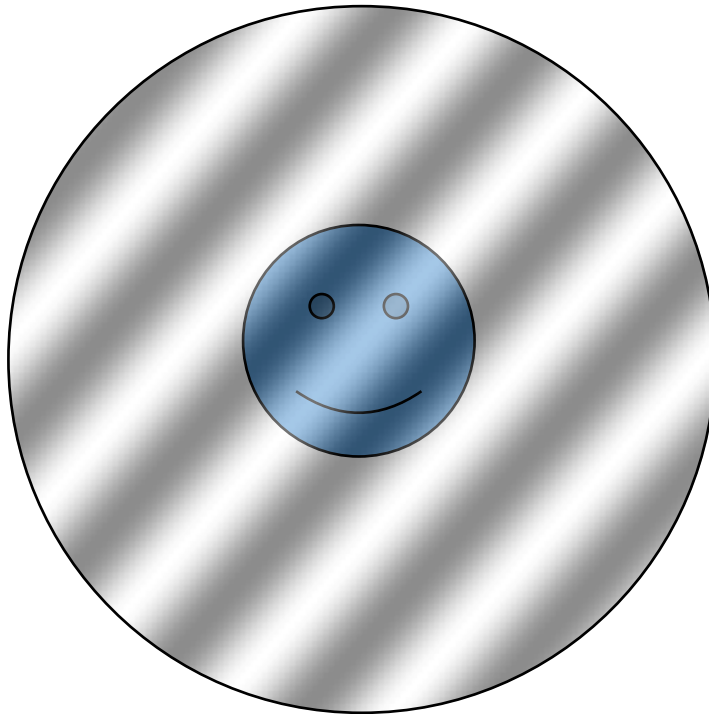


extended source

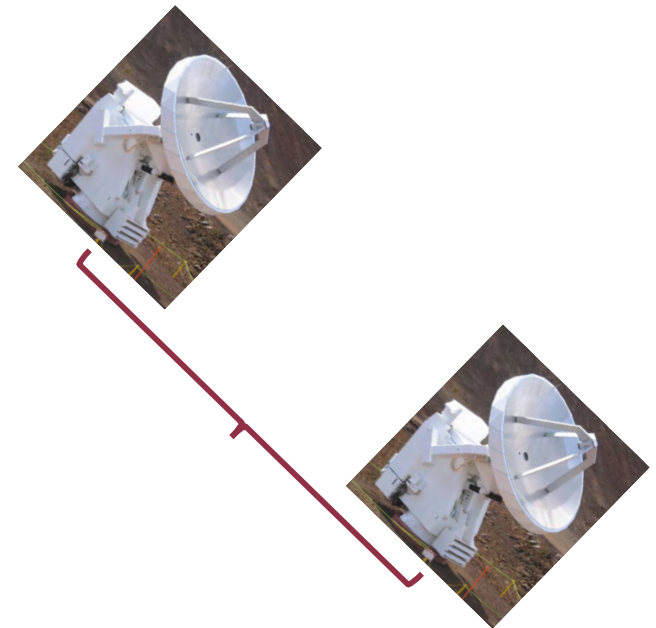


long baseline
narrow fringe pattern

$$V(u, v) = \iint T(l, m) e^{-i2\pi(ul+vm)} dl dm$$



extended source



long baseline
narrow fringe pattern
different orientation

Aperture Synthesis

basic idea: sample $V(u,v)$ at enough (u,v) points using distributed small apertures to synthesize a large aperture of size (u_{max}, v_{max})

use **more antennas** for more samples

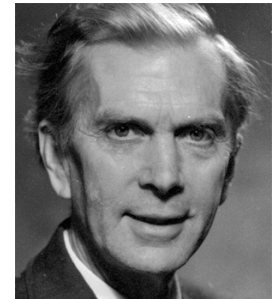
- one pair of antennas = one baseline
= two (u,v) samples at one time
- N antennas = $N(N-1)$ samples at one time
- reconfigure physical layout of N antennas for more samples (as long as source structure doesn't change with time)

use **Earth rotation** for more samples

- baseline length/orientation relative to sky change with time

use **more wavelengths** for more (continuum) samples

- u and v are measured in *wavelengths*
- “multi-frequency synthesis”: determine structure at a fiducial wavelength together with change with wavelength

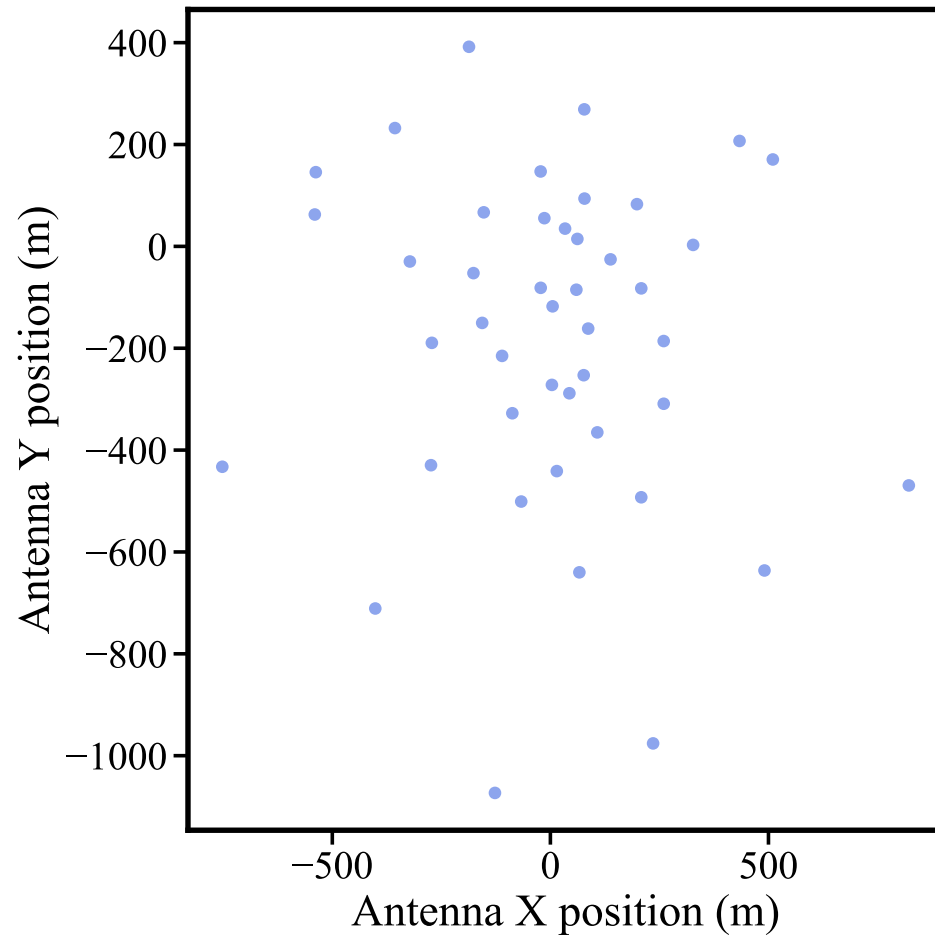


Sir Martin Ryle
1918-1984



1974 Nobel
Prize in Physics

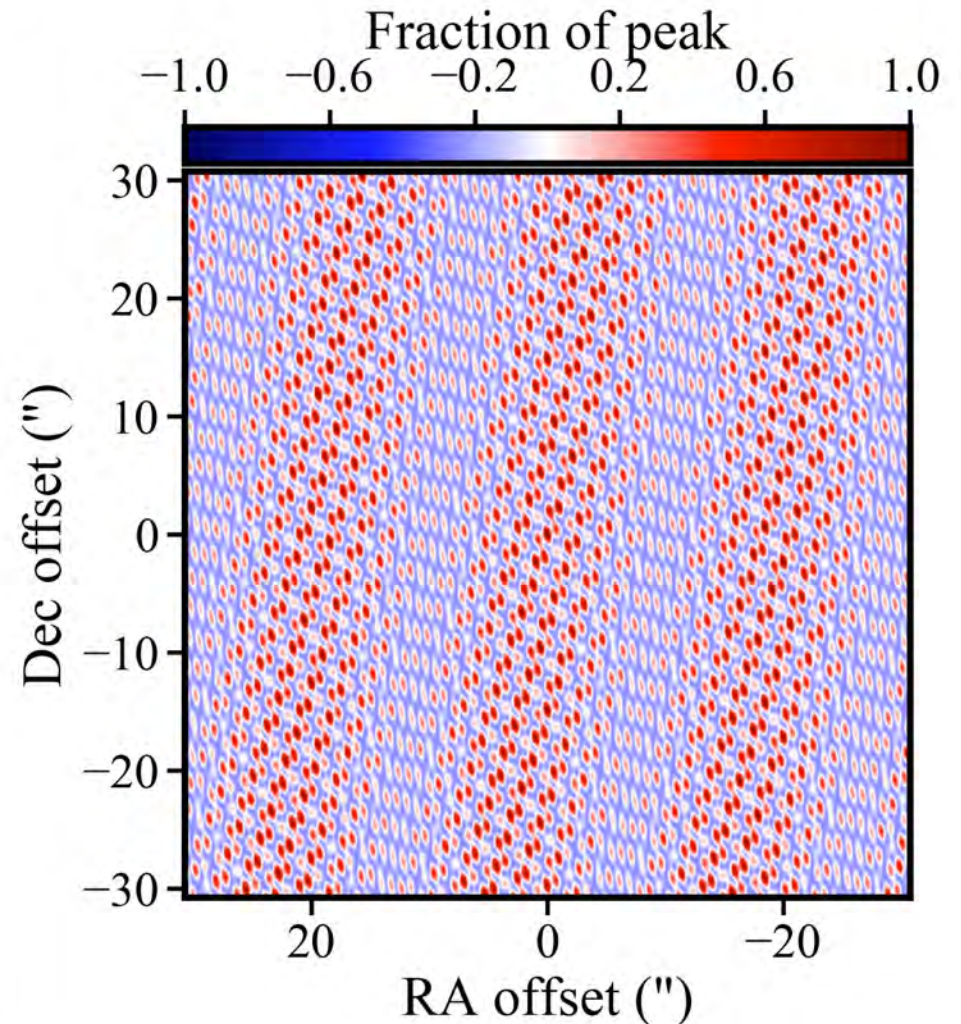
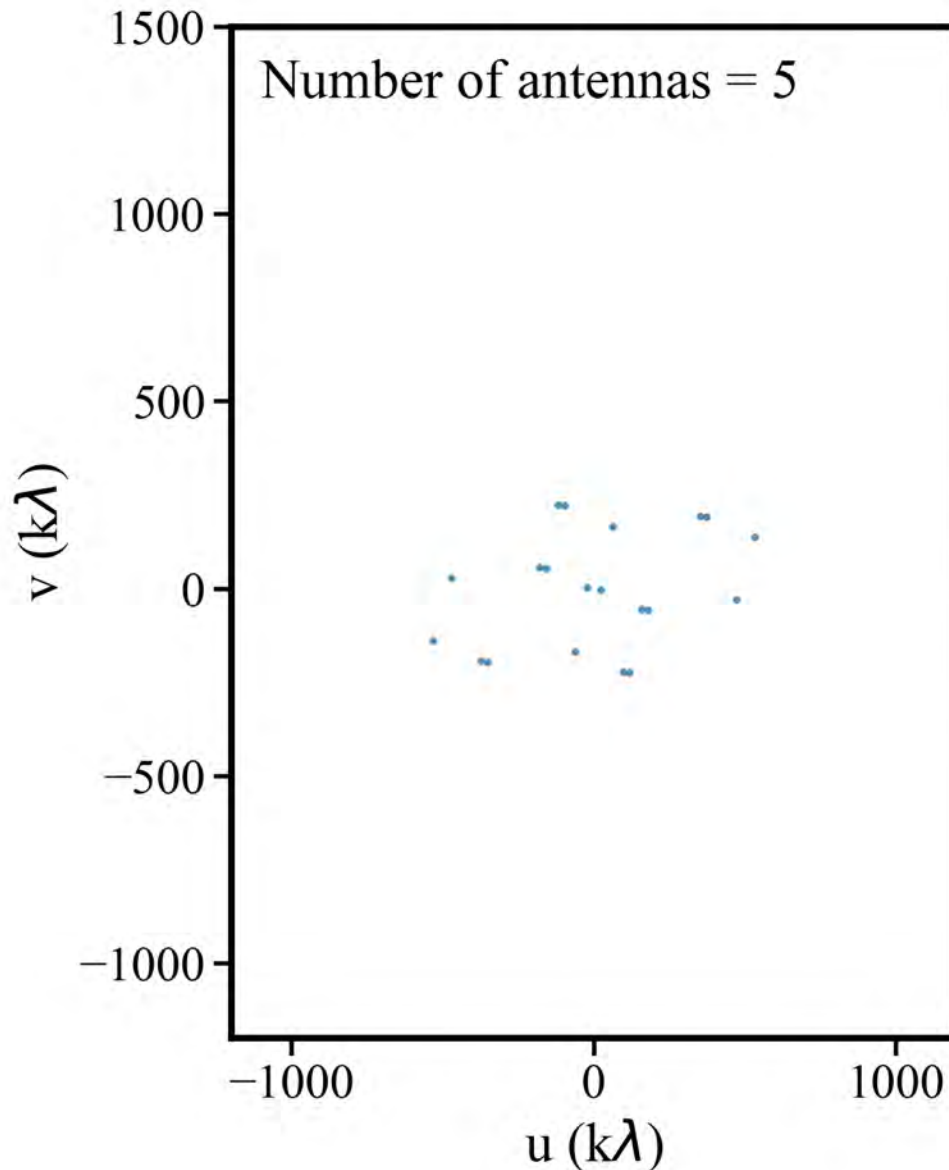
Example of (u,v) Plane Sampling



ALMA 12m antenna
locations on Aug 8, 2015

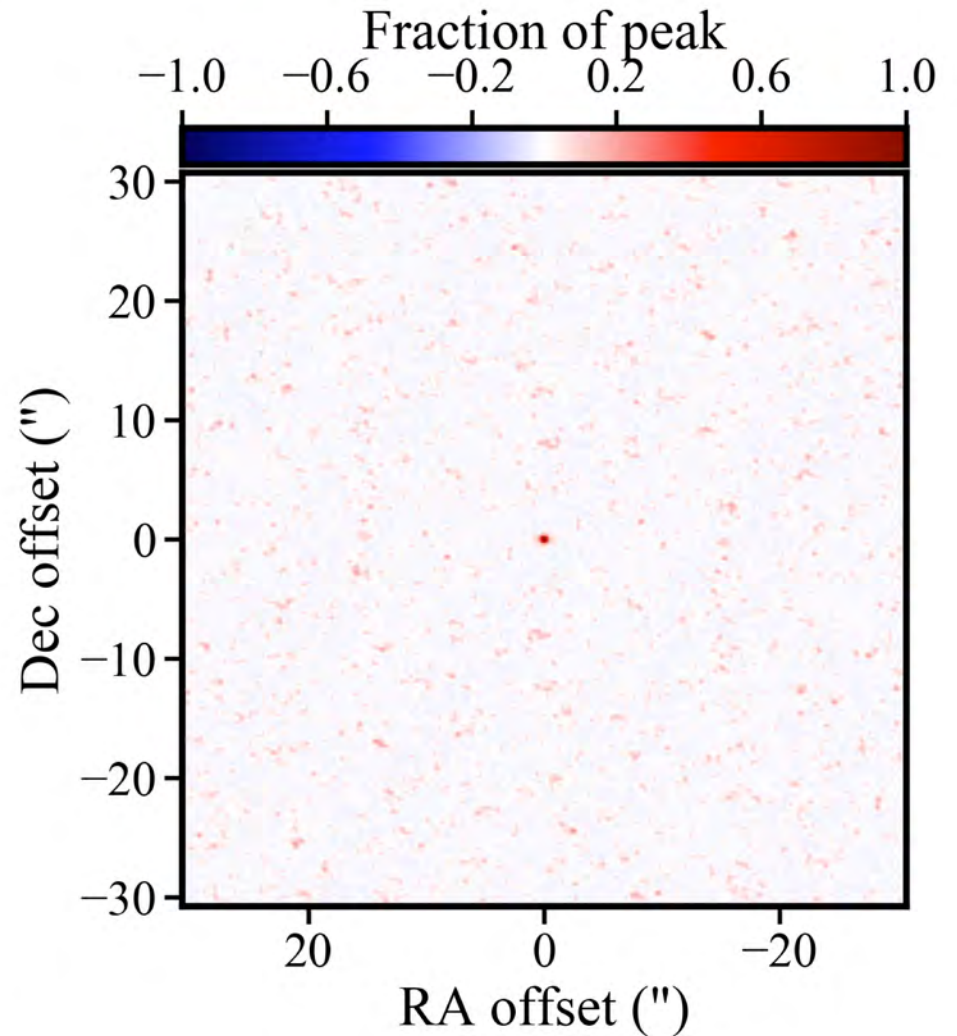
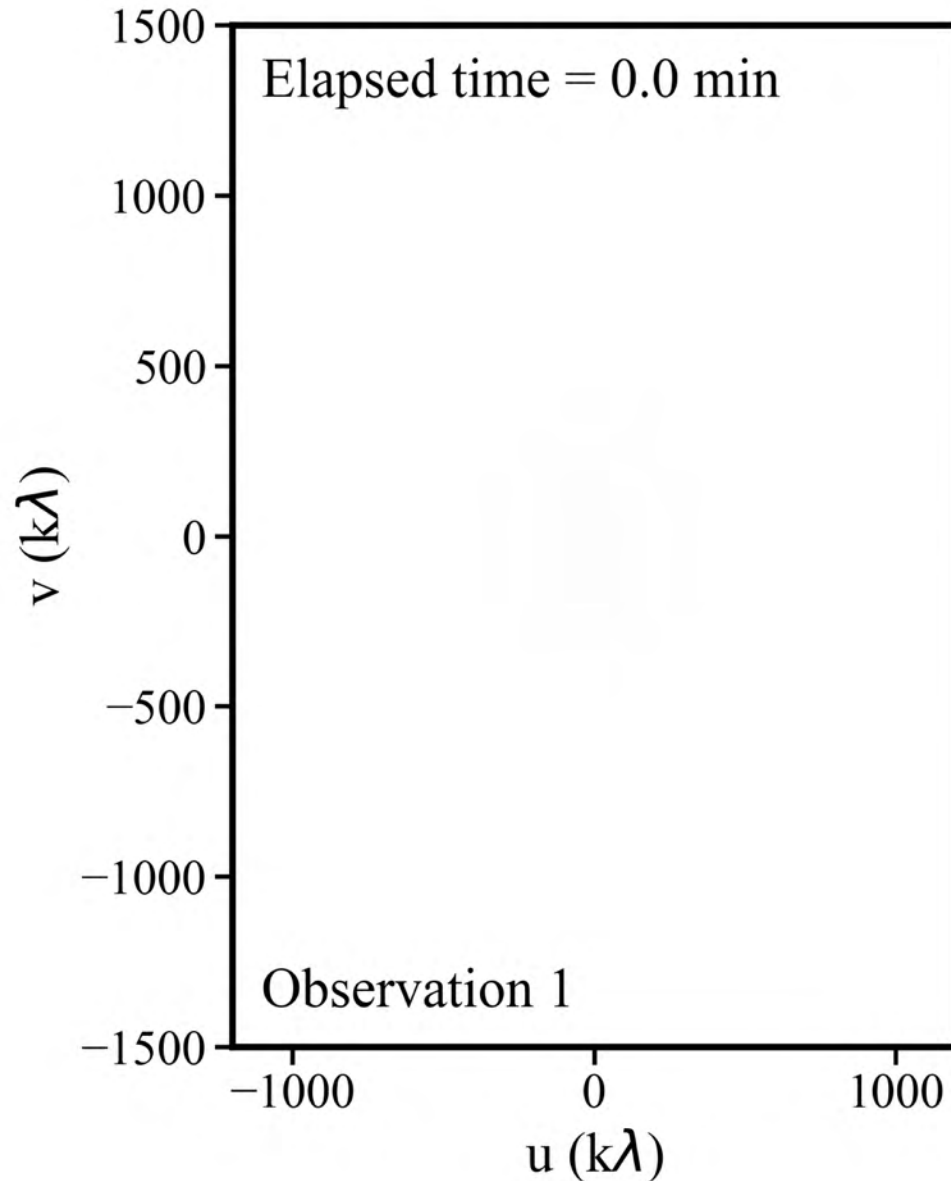
[array design: Eric Murphy lecture]

(u, v) Plane Sampling: more Antennas



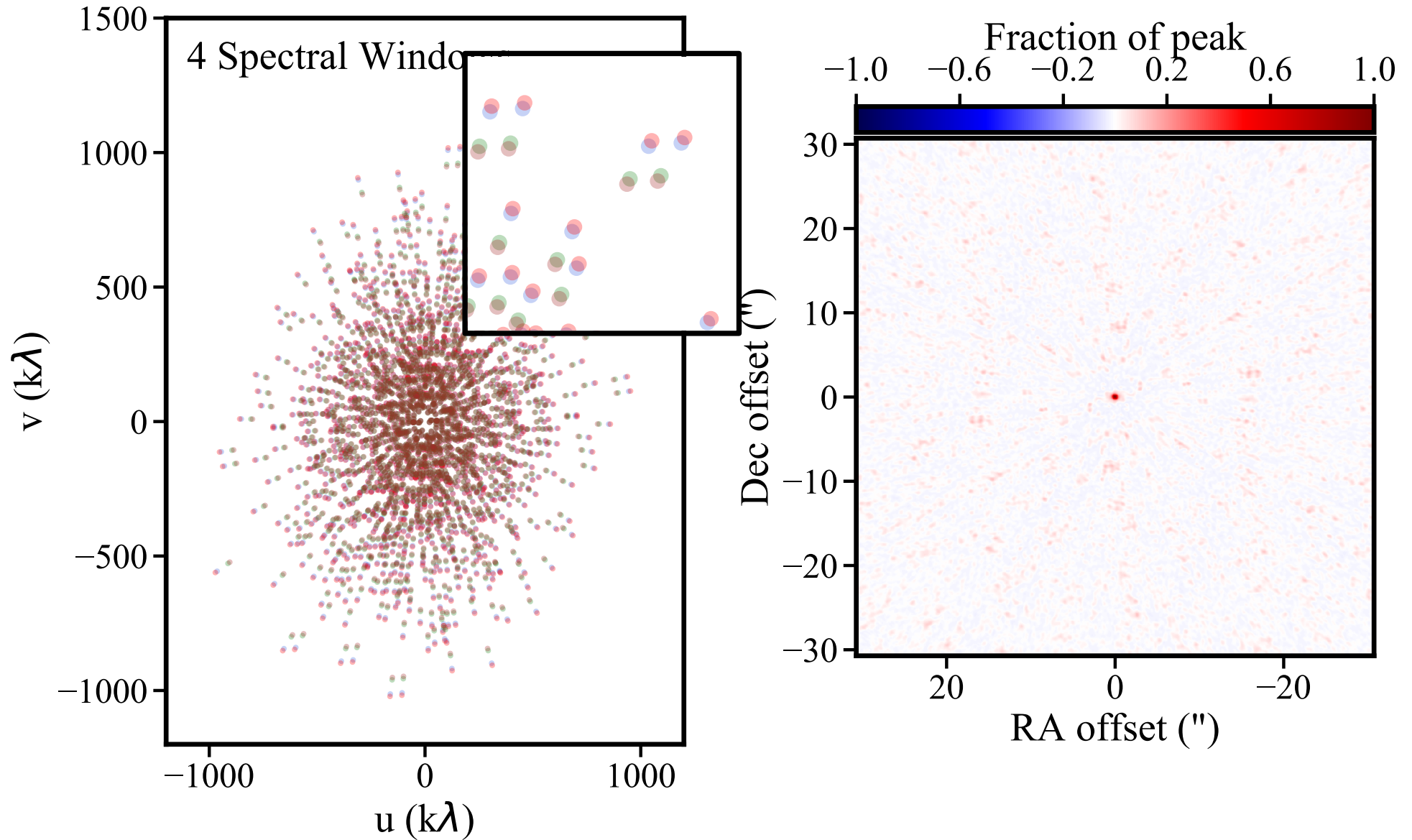
source dec -51 deg
 $\lambda = 1.3$ mm

(u, v) Plane Sampling: Earth Rotation



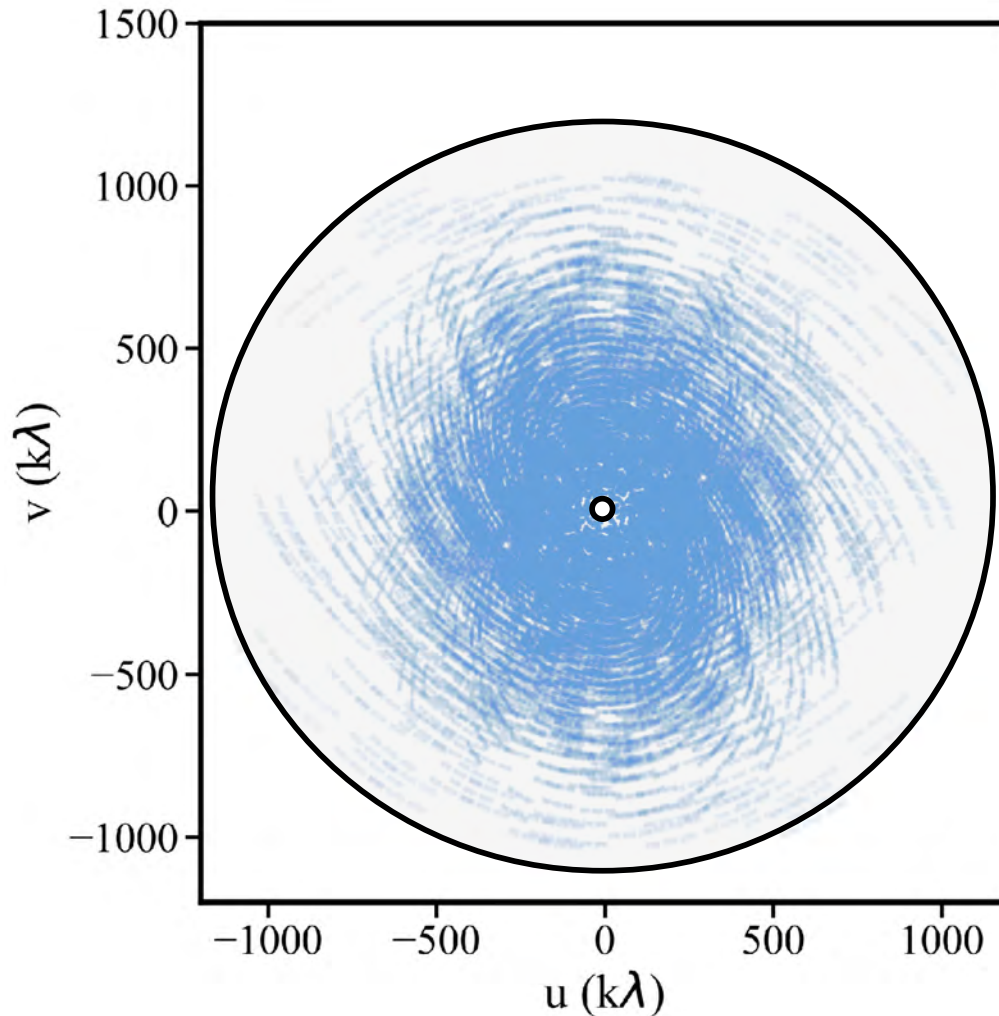
source dec -51 deg
 $\lambda = 1.3$ mm

(u,v) Plane Sampling: more wavelengths



Implications of (u,v) Plane Sampling

samples of $V(u,v)$ are limited by array and Earth-sky geometry



outer boundary

- no info on smaller scales
- resolution limit

inner boundary

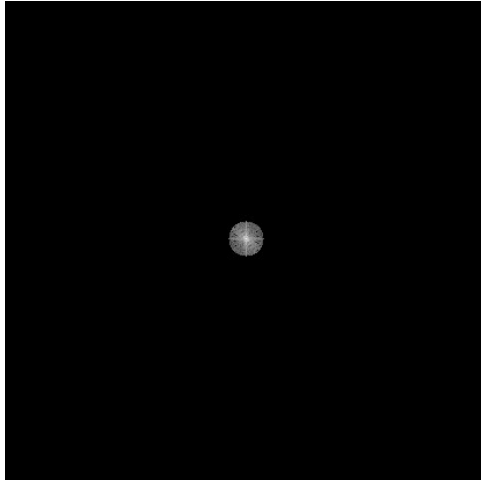
- no info on larger scales
- extended sources invisible

irregular coverage in between

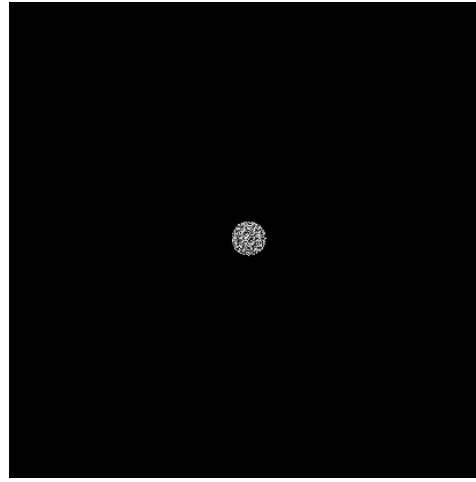
- sampling theorem violated
- information missing

Inner and Outer (u,v) Boundaries

$V(u,v)$ amplitude



$V(u,v)$ phase

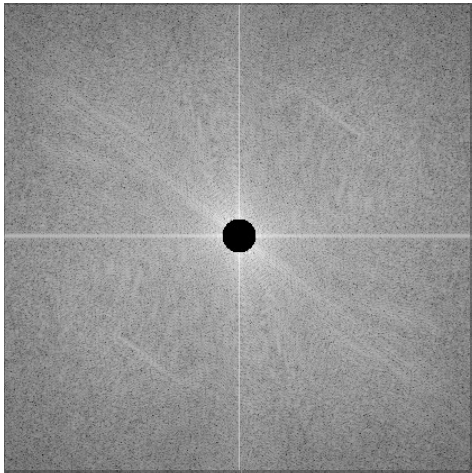


$\mathcal{F}$
 $\longrightarrow$

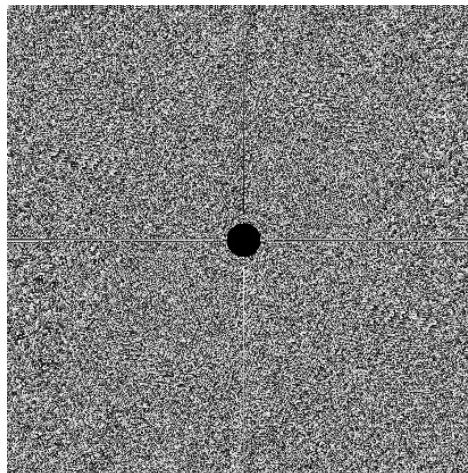
$T(l,m)$



$V(u,v)$ amplitude



$V(u,v)$ phase



$\mathcal{F}$
 $\longrightarrow$

$T(l,m)$



Calibrated Visibilities... Now what?

analyze $V(u,v)$ samples directly by model fitting [

- good for simple structures, e.g. point sources, symmetric rings
- for a purely statistical description of sky brightness, e.g. fluctuations
- visibilities have well defined noise properties

[Dom Pesce lecture]

recover an image from incomplete, noisy samples of $V(u,v)$

- Fourier transform $V(u,v)$ to create a distorted $T^D(l,m)$
- account for incomplete sampling to create a model of $T(l,m)$
- work with the model of $T(l,m)$ to do science

CASA
tclean

Formal Description of Imaging

- sample Fourier domain at discrete points

$$S(u, v) = \sum_{k=1}^M \delta(u - u_k, v - v_k)$$

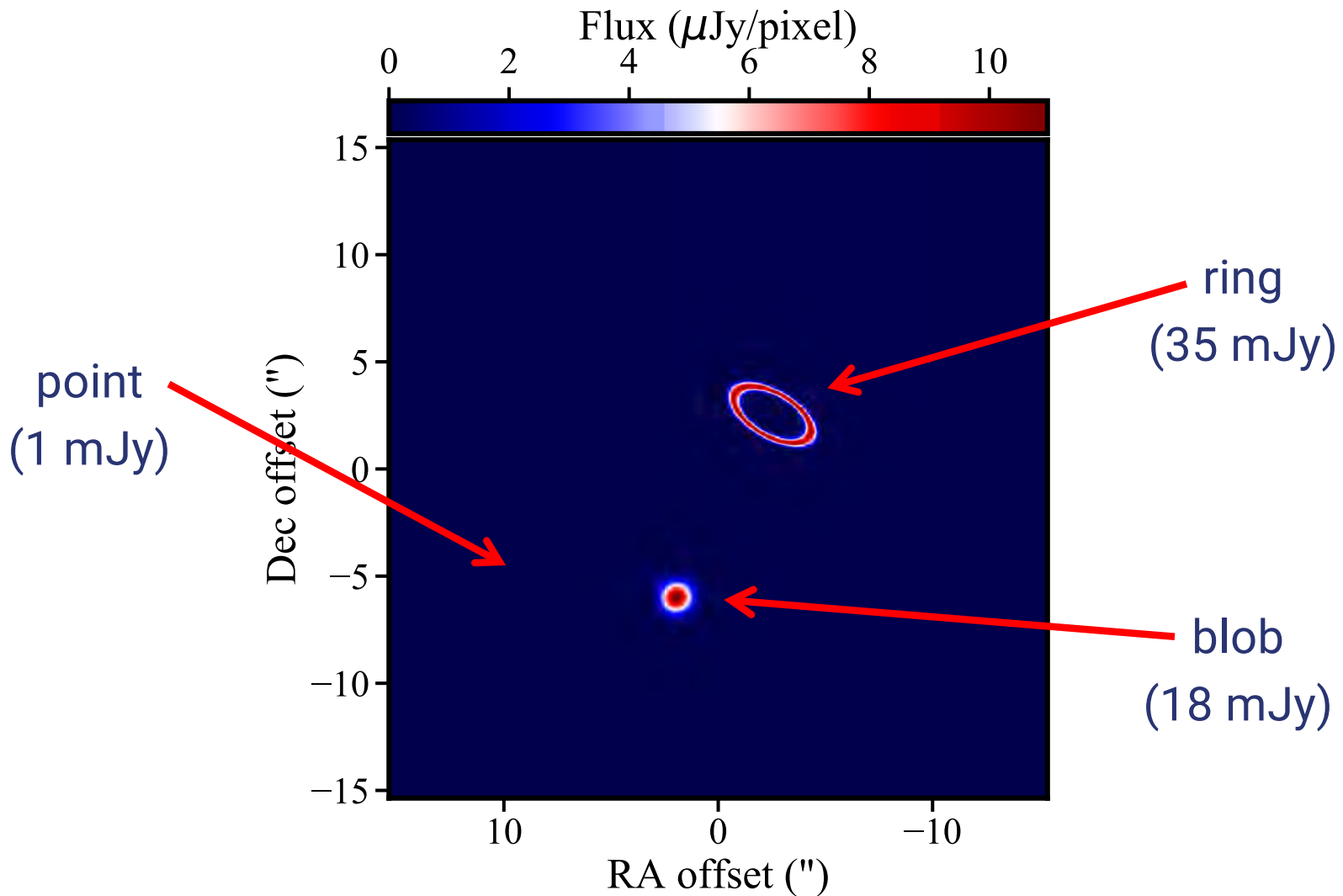
- Fourier transform sampled visibilities $V(u, v)S(u, v) \xrightarrow{\mathcal{F}} T^D(l, m)$

- apply the convolution theorem $T(l, m) * s(l, m) = T^D(l, m)$

where the Fourier transform of the sampling pattern $s(l, m) \xrightarrow{\mathcal{F}} S(u, v)$ is the “point spread function” or the “synthesized beam”

- radio astronomy jargon:
the “dirty image” is the true image convolved with the “dirty beam”

Example model sky brightness $T(l,m)$



Dirty Beam and Dirty Image

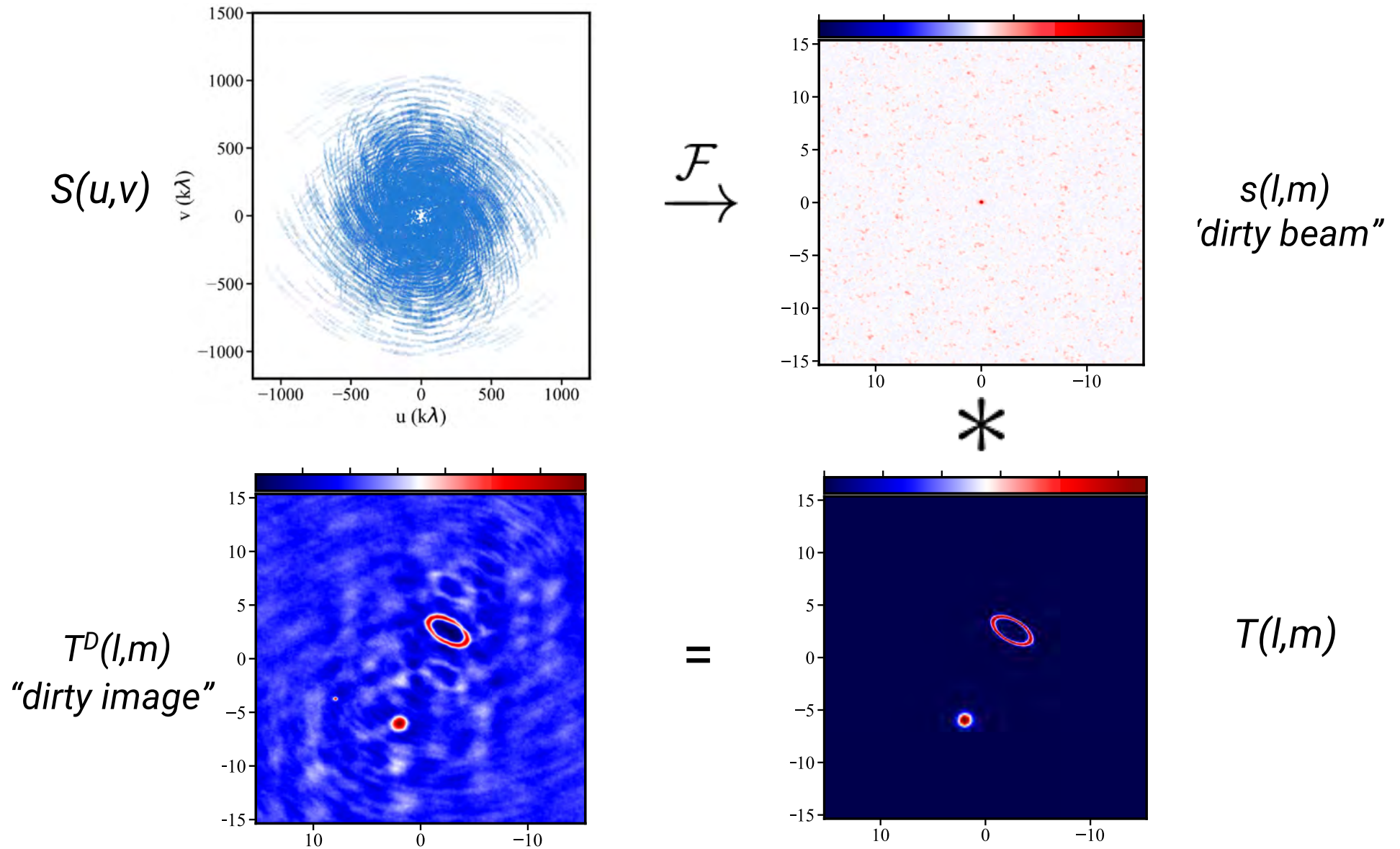


Image Size and Pixel Size

image size

- antenna response modifies the sky brightness distribution
 $T(l,m) \rightarrow T(l,m)A(l,m)$
where $A(l,m)$ is antenna “primary beam” (fwhm $\sim \lambda/D$) [Rob Selena lecture]
- natural choice for image size is often full extent of primary beam
e.g. ALMA 12 m antennas at 1.3 mm $\rightarrow$ image size 2 x 27 arcsec
- “primary beam correction”
divide final image by $A(l,m)$ to account for $A(l,m)$ attenuation;
raises noise away from the center

pixel size

- satisfy sampling theorem for longest baselines: $\Delta l < \frac{1}{2u_{max}}$ $\Delta m < \frac{1}{2v_{max}}$
- in practice, 3 to 5 pixels across dirty beam main lobe to aid deconvolution
e.g. ALMA at 1.3 mm, baselines to 1 km $\rightarrow$ pixel size < 0.05 arcsec

“Fourier Transform”

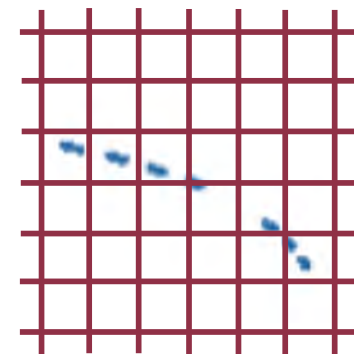
- direct computation of Fourier integral by summation is slow, generally not practical for modern imaging applications
- FFT = Fast Fourier Transform algorithm
- trace to Gauss' work in 1805 to interpolate orbit of asteroids Pallas and Juno; (re)discovered by Cooley & Tukey 1965
- $O(N^2 \log N)$ vs. $O(N^4)$ for N^2 image cells
- optimized codes readily available
- **FFT requires data on a regularly spaced grid**
- but aperture synthesis does not provide $V(u,v)$ samples on a regularly spaced grid, so...

Gridding

Gridding is used to resample $V(u,v)$ onto a regular (u,v) grid for FFT

- conventional approach is to use convolution
- (u,v) cell size $\approx 0.5 \times D$, where D = antenna diameter

$$V^G(u, v) = V(u, v)S(u, v) * G(u, v) \\ \xrightarrow{F} T^D(l, m)g(l, m)$$



- other gridding steps may include
 - functions to apply primary beam weighting and offsets (“mosaicking”)
 - functions to apply wide-field phase shifts (“W projection”)
 - functions to correct for primary beam differences (“A projection”)

[Brian Mason, Urvashi Rao, Preshanth Jagannathan lectures]

Visibility Weighting Schemes

Introduce weighting function $W(u,v)$ into the gridding process

- $W(u,v)$ modifies the sampling function: $S(u,v) \rightarrow S(u,v)W(u,v)$
- changes the dirty beam shape, $s(l,m)$

Natural weighting: $W(u,v) = 1/\sigma^2(u,v)$ where σ^2 is noise variance

- best point source sensitivity, lower resolution, more dirty beam structure

Uniform weighting: $W(u,v) = 1/\delta(u,v)$ where δ is density of (u,v) points

- higher noise, higher resolution, less dirty beam structure

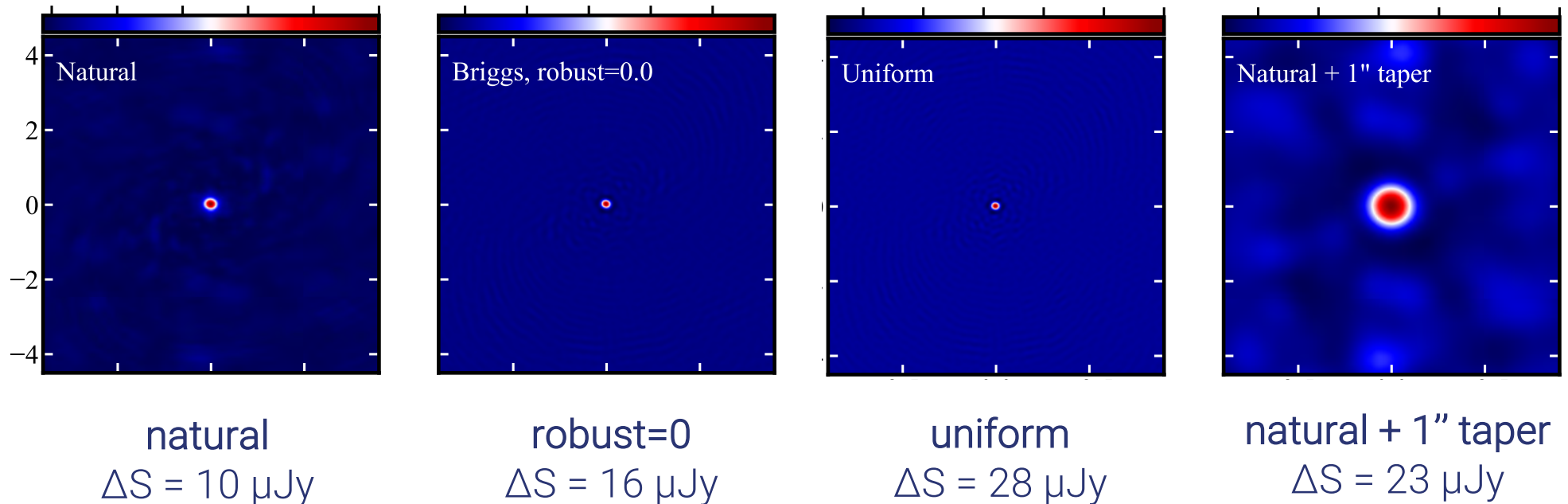
Robust (“Briggs”) weighting: $W(u,v)$ combines noise and density

- adjustable parameter allows for continuous variation between natural and uniform weighting (in CASA from -2 to 2)
- usually can obtain most of natural weight sensitivity at the same time as most of uniform weight resolution (!)

Tapering: apodize (u,v) sampling by a Gaussian function

- lower resolution to improve sensitivity to extended structure

Example Dirty Beams (same data!)



- imaging parameters provide a lot of freedom
- appropriate choice depends on science goal, e.g.
 - point source detection: natural weight
 - fine detail of strong source: uniform weight
 - weak and extended emission: taper

Beyond Dirty Images: Deconvolution

- to keep you awake at night:
 - ∃ an infinite number of $T(l,m)$ compatible with sampled $V(u,v)$

Deconvolution

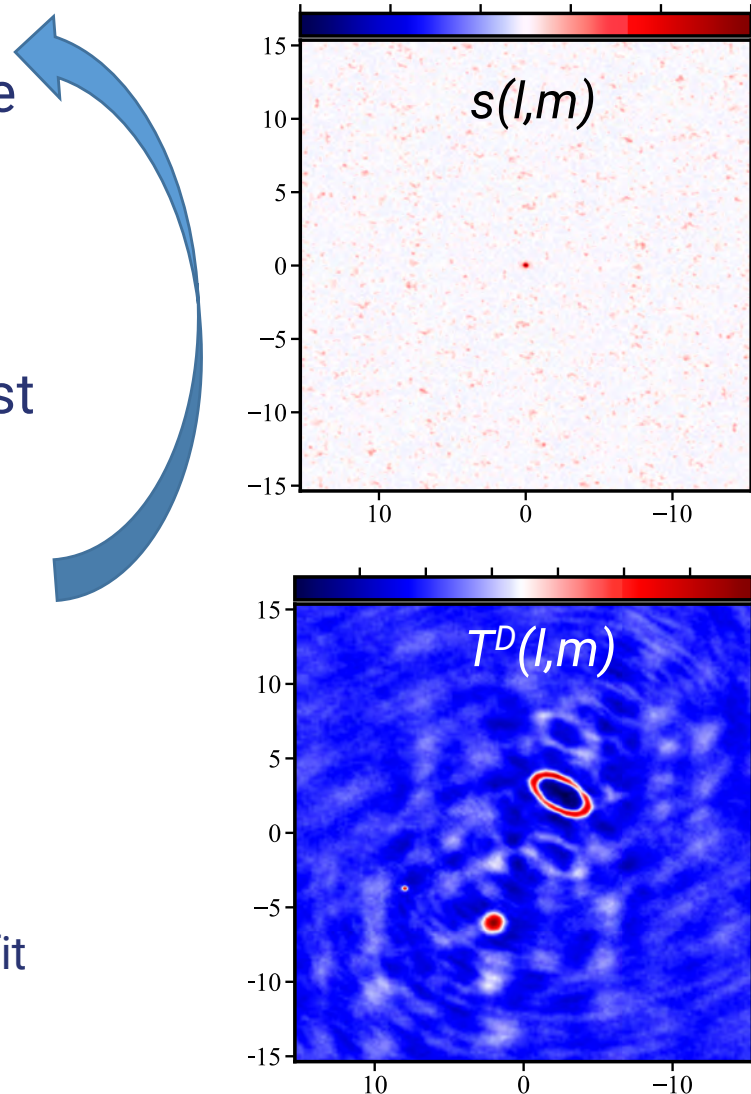
- use non-linear techniques to interpolate/extrapolate $V(u,v)$ samples into unsampled regions of (u,v) plane to find a *plausible* model of $T(l,m)$
- requires *a priori* assumptions about $T(l,m)$ to fill unsampled (u,v) plane
- dominant deconvolution algorithm in radio astronomy is “clean”
 - *a priori* assumption: $T(l,m)$ can be represented by point sources
 - many variants have been developed for computational efficiency and performance on extended structure (Clark, Cotton-Schwab, Multi-Scale, ...)
- a very active research area!
 - see EHT paper for some examples, e.g. regularized likelihood methods

Classic Hogbom Clean Algorithm

1. find highest peak of residual image
2. subtract scaled dirty beam $s(l,m)$ x “loop gain” from this peak
3. add this point source position and amplitude to a clean component list and to a model image
4. if peak of residual image > stopping threshold, then goto 1

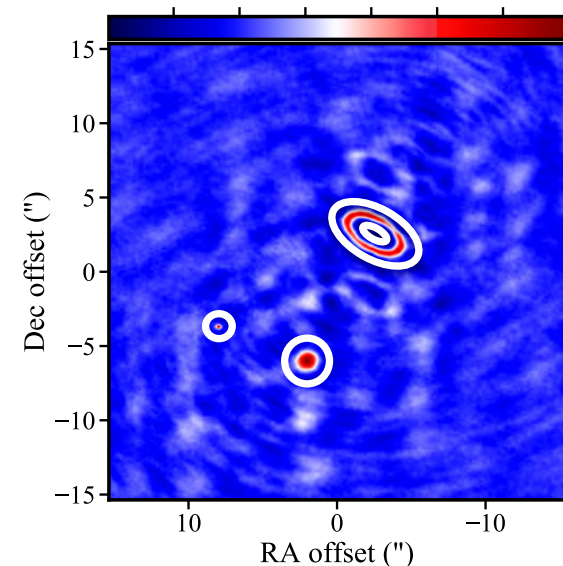
create final **restored image**

- a) convolve the model image by a “clean beam” = an elliptical Gaussian fit to the main lobe of the dirty beam
- b) add back residual image = noise + residual source structure

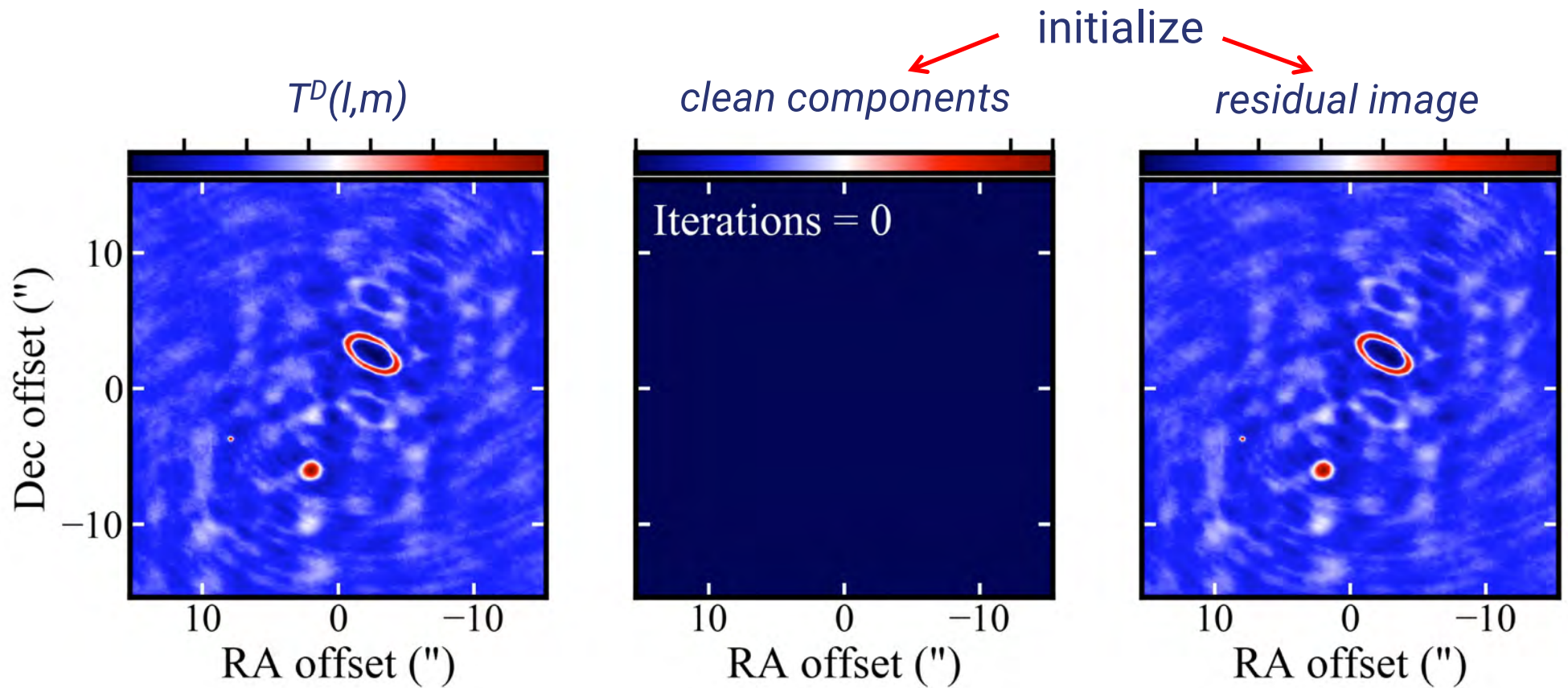


Clean Algorithm Parameters

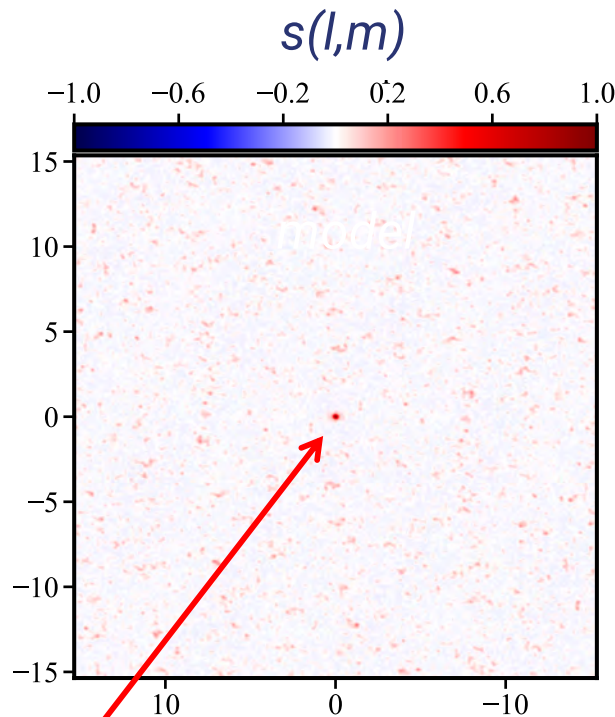
- loop gain
 - good results for 0.1 – 0.3 (CASA tclean default = 0.1)
 - lower values can help with smooth and extended emission
- stopping threshold
 - peak < threshold = multiple of theoretical rms noise, e.g. 3 x rms
 - peak < threshold = fraction of dirty map maximum (useful if strong sources prevent theoretical rms threshold from being reached)
- finite support (image masks)
 - include *a priori* information about where to search for clean components in image
 - useful, often essential, for best results
 - can be dangerous
 - can be an arduous manual process, but modern automasking algorithms work well



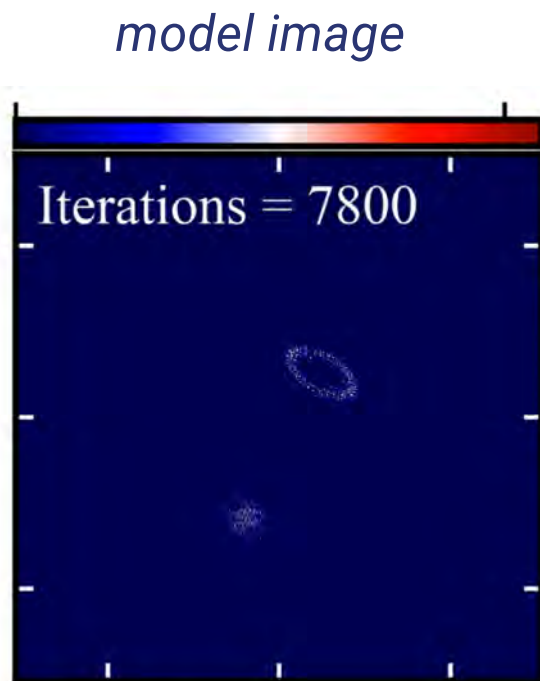
Clean Example



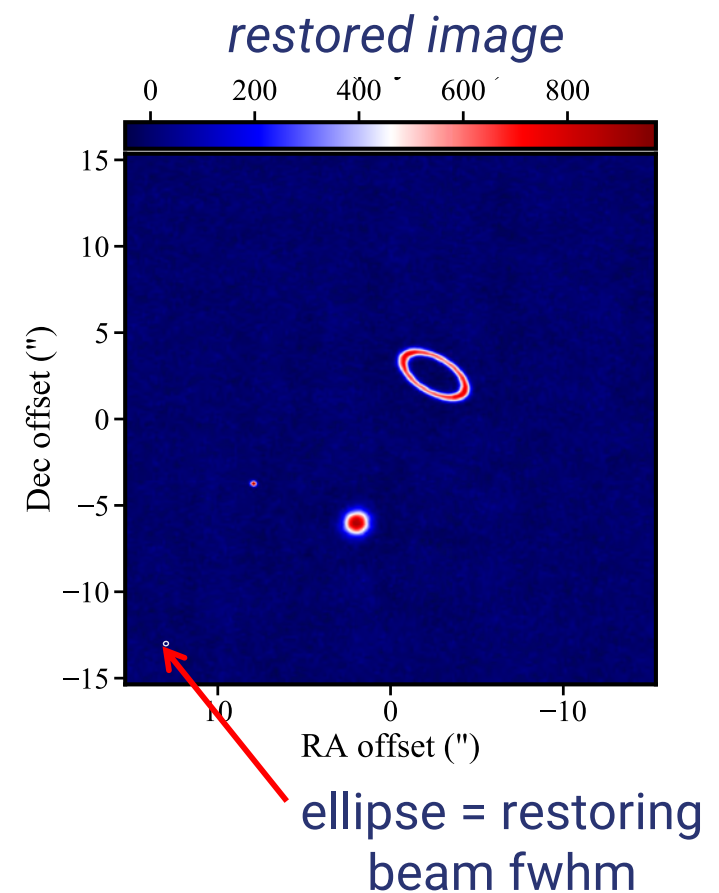
Clean Example: Restored Image



fit Gaussian to $s(l,m)$



convolve model image
by fitted Gaussian and
add back residuals

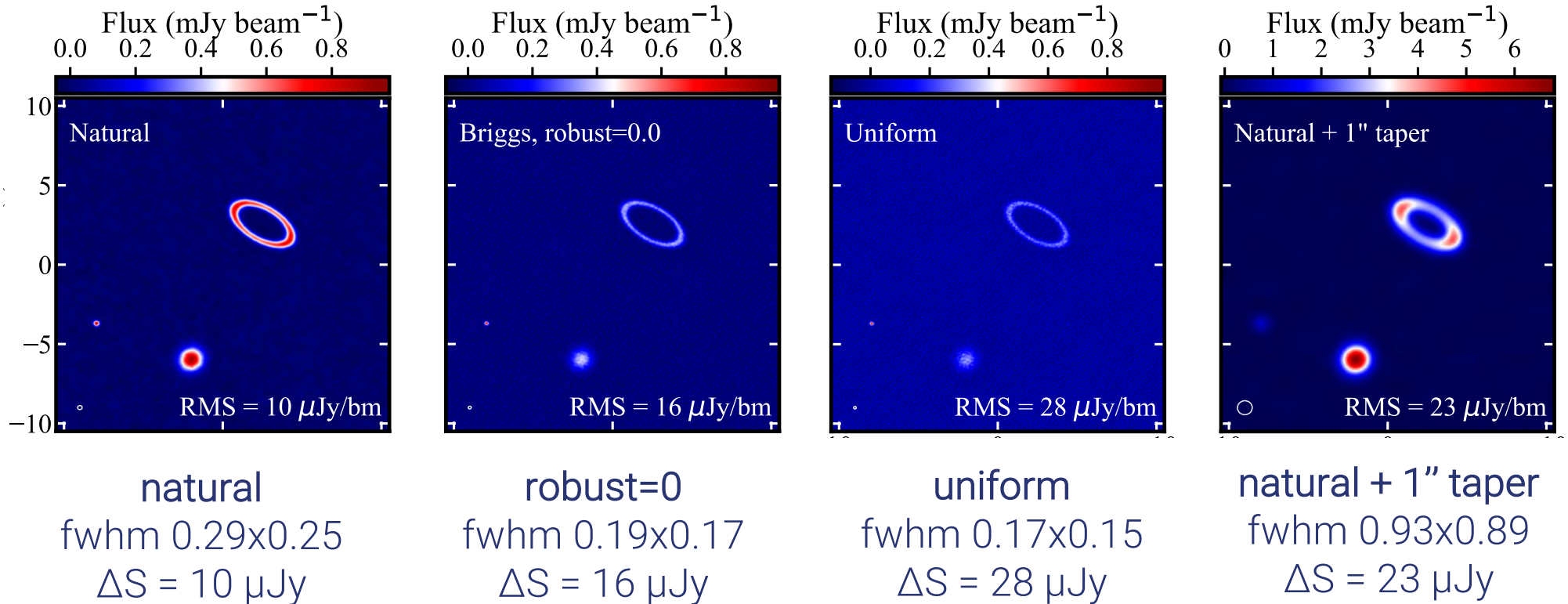


final image depends on

- imaging parameters: pixel size, visibility weighting scheme, gridding, ...
- deconvolution: algorithm, parameters, iterations, ...

Clean: Different Weighting Schemes

- restored images emphasize different angular scales from the $V(u,v)$ samples



CASA `tclean` filename extensions

- `<imagename>.image`
 - restored image
- `<imagename>.psf`
 - point spread function (dirty beam)
- `<imagename>.model`
 - model image after deconvolution, i.e. clean components
- `<imagename>.residual`
 - residual image, i.e. after subtracting clean components
- `<imagename>.mask`
 - deconvolution mask
- `<imagename>.pb`
 - primary beam model
- `<imagename>.pbcor`
 - primary beam corrected image

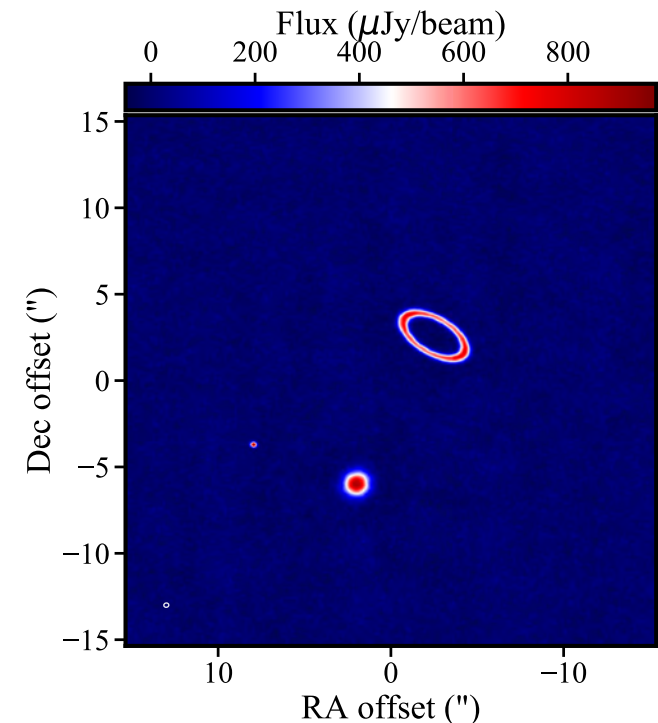
Measures of Image Quality

dynamic range

- ratio of peak brightness in image to rms noise in region devoid of emission
- easy way to calculate a *lower limit* to the error in brightness in a non-empty region

e.g. peak 880 and rms 10 $\mu\text{Jy}/\text{beam}$

→ dynamic range = 88



fidelity

- difference between reconstructed image and the true image
- fidelity = input model/difference = inverse of the relative error
= $\text{model} * \text{beam} / \text{abs}(\text{model} * \text{beam} - \text{reconstructed image})$
- generally much lower than implied by dynamic range

Spectral Lines and Polarization

- discussion so far applies to a single spectral channel of width $\delta\nu$
- science may require many such channels across total bandwidth $\Delta\nu$
 - spectral lines from molecular or atomic transitions
 - a significant continuum spectral slope or curvature
 - also helpful from a technical perspective
 - edit interference, avoid “bandwidth smearing” (radial smearing in image from averaging over $\Delta\nu$ that limits useable field of view)
 - each channel is imaged (and deconvolved) independently

[Ylva Pihlstrom lecture]

- science may require full Stokes parameters: I, Q, U, V
 - each Stokes parameter imaged (and deconvolved) independently
 - then combined to form, e.g. fractional linear polarization $\sqrt{(Q^2+U^2)} / I$

[Frank Shinzel lecture]

Concluding Remarks

- interferometry samples Fourier components of sky brightness
- make images by Fourier transforming the sampled visibilities
- deconvolution attempts to correct for incomplete sampling
- remember
 - there are an infinite number of images compatible with the visibilities
 - missing (or corrupted) visibilities affect the entire image
 - astronomers must make decisions in imaging and deconvolution
- it is fun and worth the trouble → high resolution images!

many, many issues not covered in this talk,

see lectures on advanced imaging techniques and references