

Imaging a Pulsar's Scattering Disk

Happy 70th Miller!

Walter Briskeen

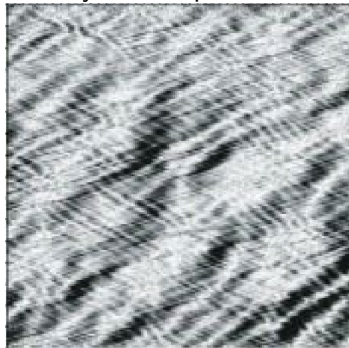
National Radio Astronomy Observatory

2011 May 19



Scintillation Arcs

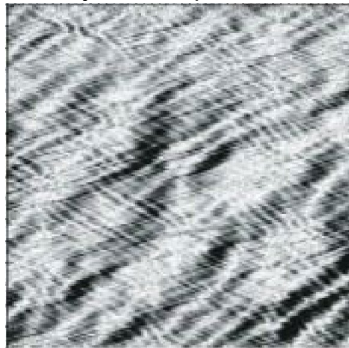
Dynamic Spectrum



$$I(\nu, t)$$

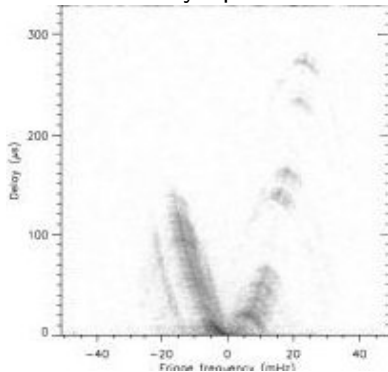
Discovery of Scintillation Arcs

Dynamic Spectrum



$$I(\nu, t)$$

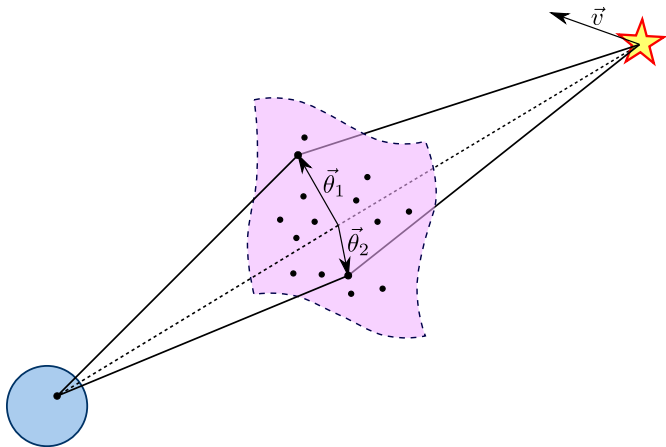
Secondary Spectrum



$$S(T, R) \equiv |\mathcal{F}[I(\nu, t)]|^2$$

see *Stinebring et al., 2001 and Walker et al., 2005*

Diffractive Scintillation Thin Screen Geometry



- * Distance to pulsar $\equiv D_{\text{psr}}$
- * Distance to screen $\equiv D_{\text{scr}} = (1 - s)D_{\text{psr}}$
- * Effective distance $\equiv D_{\text{eff}} = \frac{1-s}{s}D_{\text{psr}}$
- * Effective velocity $\equiv \vec{V}_{\text{eff}} = \frac{1-s}{s}\vec{v}$

see Cordes et. al, 2006

Model of Parabolic Arcs

Delay coordinate

- * From geometry

$$T \equiv \tau_1 - \tau_2 = \frac{D_{\text{eff}}}{2c} (\theta_1^2 - \theta_2^2)$$

Doppler rate coordinate

- * From time derivative of τ :

$$R = \vec{V}_{\text{eff}} \cdot (\vec{\theta}_1 - \vec{\theta}_2)$$

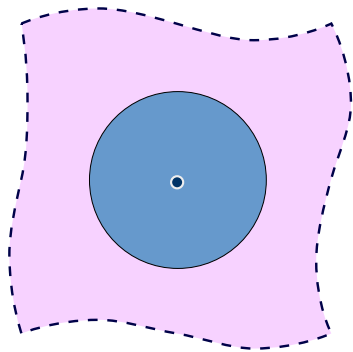
The parabola

- * Assume dominating central concentration near $\vec{\theta}_2 = 0$
- * Then:

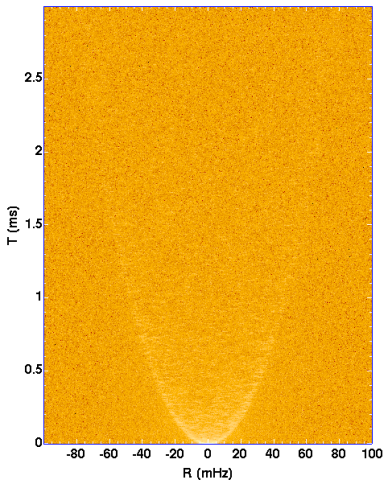
$$T \geq \frac{\lambda^2 D_{\text{eff}}}{2c V_{\text{eff}}^2} R^2$$

- * Equality occurs for $\vec{\theta}_1 \parallel \vec{V}_{\text{eff}}$

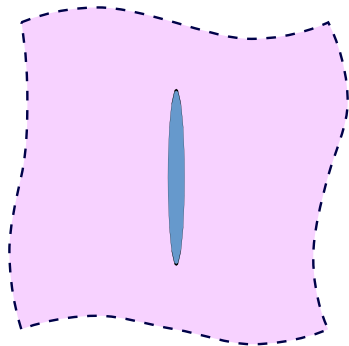
Example 1



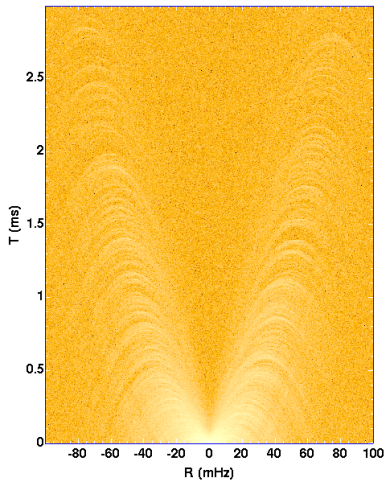
Core Disk Distribution



Example 2



10:1 Axial ratio \parallel to velocity



VLBI Study of Scattering Screen

Goal 1

- * Investigate arcs at very high frequency resolution

Goal 2

- * Validate parabolic arc model
- * Make model-independent image of scattering screen

Goal 3

- * Break degeneracies
 - Measure anisotropy of scattering
 - Determine orientation for improved interpretation of single-dish data

The Ad-hoc VLBI Array

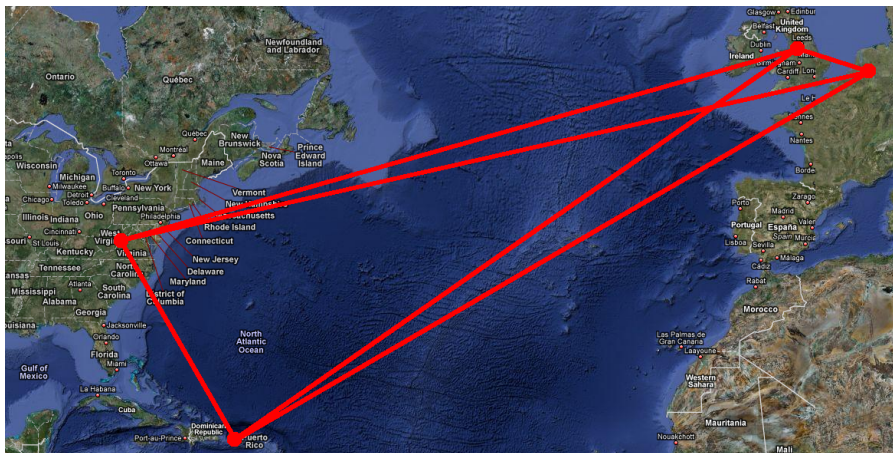


Image courtesy of Google

- * Need: Low freq, long baselines, high sensitivity & mutual visibility
- * Array: GB (100 m), AO (305 m), JB (76 m) and WB (93 m equiv.)

Experimental Setup

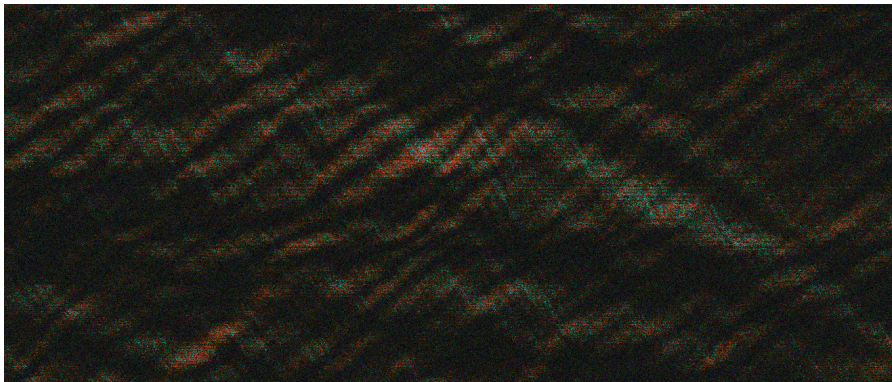
The observation

- * Target: pulsar B0834+06
- * 2 hours on source
- * Frequency: 310 to 342 MHz with dual circular polarization

Correlation

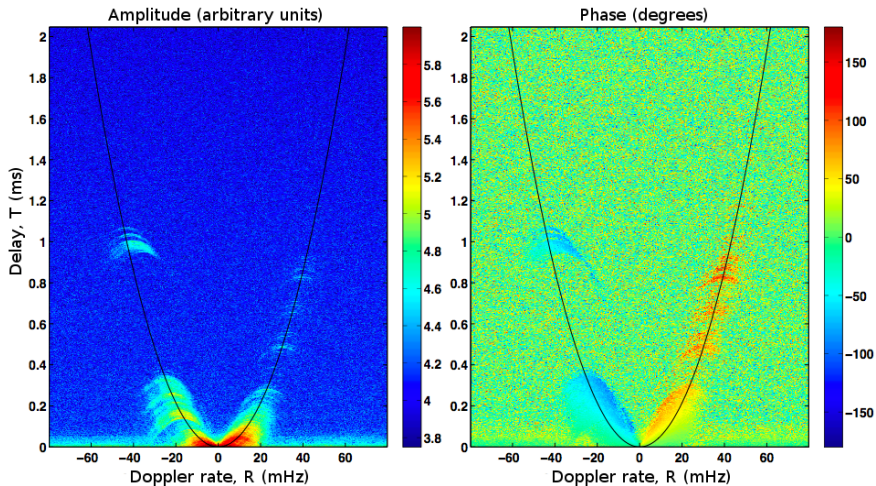
- * Used Adam Deller's nascent DiFX software correlator at Swinburne Univ.
- * 131072 spectral channels (244 Hz resolution)
- * 1.25 second integrations
- * Pulsar gate used to boost signal-to-noise ratio

Visibility Dynamic Spectrum (AR-GB)



- * ~ 600 seconds \uparrow of data over ~ 200 kHz \rightarrow
- * Amplitude mapped to intensity
- * Phase mapped to color (red to blue)
- * Note, only a very small piece of the dynamic spectrum is shown

Visibility Secondary Spectrum (AR-GB)



see *Briskin et al. 2010*

Mining a rich data set

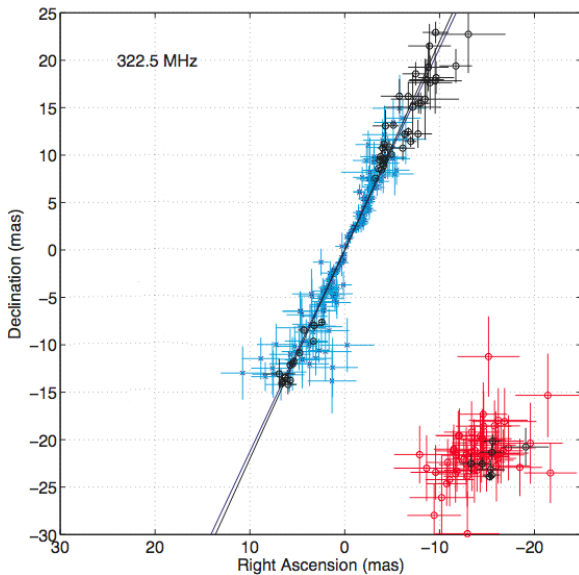
For each point in the secondary spectrum

- * Delay, T
- * Doppler rate, R
- * Amplitude (on each baseline)
- * Phase (on each baseline)

Phases yield sky coordinates

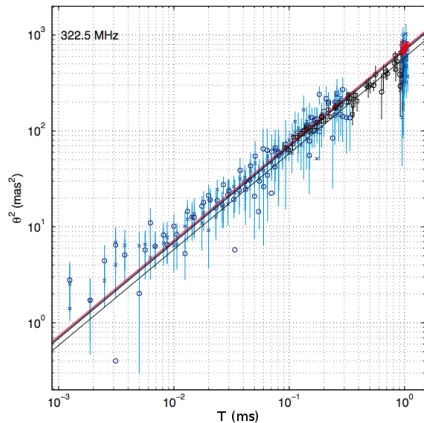
- * $\phi = \frac{2\pi\nu}{c}(ul + vm)$ (as for interferometry)
- * Can directly test $T = \frac{D_{\text{eff}}}{2c}\theta^2$ and $R = \vec{V}_{\text{eff}} \cdot \vec{\theta} \dots$
- * ... and determine D_{eff} and \vec{V}_{eff}

Astrometrically Recovered Image

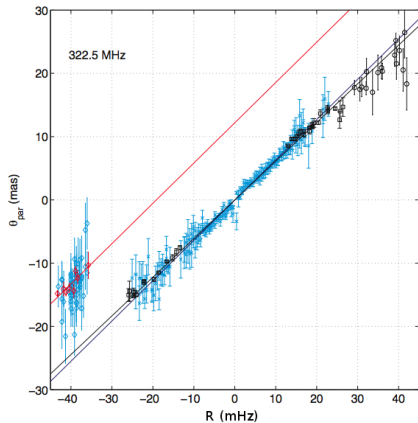


Physical Parameter Estimation

Distance fit



Velocity fit

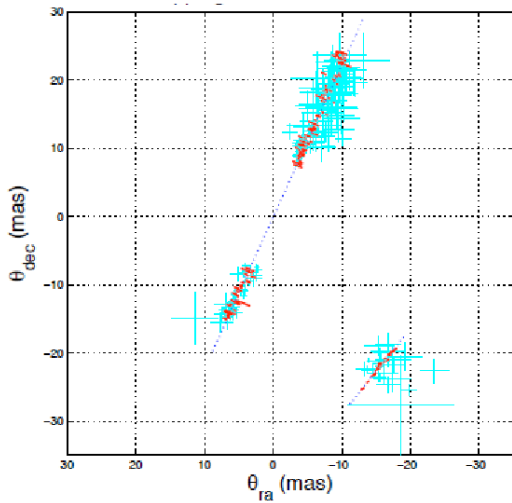
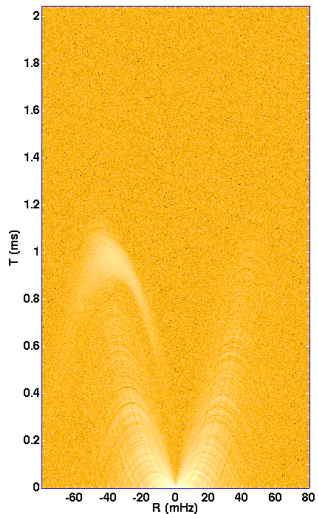


* $D_{\text{eff}} = 1171 \pm 10 \text{ pc} \rightarrow D_{\text{scr}} = 415 \pm 5 \text{ pc}$

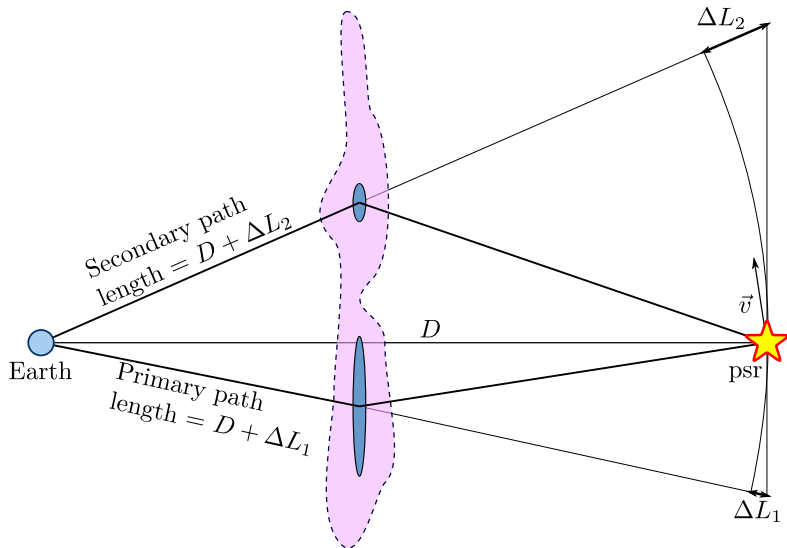
* $V_{\text{eff},\parallel} = 305 \pm 3 \text{ km s}^{-1}$

* $V_{\text{eff},\perp} = -150 \pm 5 \text{ km s}^{-1}$

Model Recovered Image



Pulsar timing effects



Conclusion

Stationary Phase Points I

Delay model

$$\tau(\vec{\theta}) = \frac{D_{\text{eff}}}{2c} \theta^2 \text{ with direct path: } \tau(\vec{\theta} = 0) \equiv 0$$

Propagation via Fresnel-Kirchhoff integral

$$\vec{E}(\nu) \propto \nu \int e^{-2\pi i \nu \tau(\vec{\theta})} d\vec{\theta} \vec{E}_{\text{psr}}(\nu)$$

In diffractive scintillation, this integral is dominated by points where constructive interference gives rise to high magnification,

$$\vec{\nabla} \tau(\vec{\theta}) = 0,$$

which are called *stationary phase points*. Their brightness contribution is related to their magnification

$$\mu = \left[\nu \nabla^2 \tau(\vec{\theta}) \right]^{-1}$$

Stationary Phase Points II

Propagation (cont.)

The Fresnel-Kirchhoff integral

$$\vec{E}(\nu) \propto \nu \int e^{-2\pi i \nu \tau(\vec{\theta})} d\vec{\theta} \vec{E}_{\text{psr}}(\nu)$$

can then be turned into a sum over stationary phase points, $\vec{\theta}_j$:

$$\vec{E}(\nu) \propto \nu \sum_j \mu_j e^{-2\pi i \nu \tau(\vec{\theta}_j)} \vec{E}_{\text{psr}}(\nu)$$

Simplifying assumption

Geometry of screen remains fixed

$$\frac{d\vec{\theta}_j}{dt} = 0$$