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18-cm VLA observations of OH towards the Galactic Centre

Absorption and emission in the four ground-state OH lines⋆

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Abstract. The OH distribution in the Sgr A Complex has been observed in the 1612-, 1665-, 1667- and 1720-MHz OH transitions with the Very Large Array (VLA) in BnA configuration. Maps are presented with a channel velocity resolution of about 9 km s^{-1} and with angular resolutions of $4'' \times 3''$. Some clear results are highlighted here, such as absorption from the Circumnuclear Disk (CND) and the OH-Streamer inside the CND near Sgr A*, strong absorption towards most of the eastern and western parts of the Sgr A East shell, lack of absorption towards both Sgr A West and the compact H II-regions to the east of Sgr A East, and double-lobed structure of the High Negative Velocity Gas (HNVG) oriented northeast and southwest of Sgr A*, and finally compact, point-like maser emission in all four transitions, in particular a 1720-MHz maser at -132 km s^{-1} in the CND as counterpart to a 1720-MHz maser at $+132$ km s^{-1} in the CND.

Key words. Galaxy: center – ISM: individual objects: Sgr A – ISM: molecules – ISM: clouds – masers – surveys

1. Introduction

The $^2\Pi_{3/2}, J = 3/2$ ground rotational state of the Hydroxyl (OH) radical is split into four hyperfine levels, leading to the well-known 18-cm main line transitions at 1665 and 1667 MHz and satellite line transitions at 1612 and 1720 MHz. The relative intensities of the lines expected from an optically thin gas in local thermodynamic equilibrium (LTE) are 1612:1665:1667:1720 = 1:5:9:1. However, the line ratios are highly sensitive to the physical environment. The first high-resolution maps of the 18-cm emission from Sgr A and the associated 1665- and 1667-MHz OH molecular gas distribution in the Galactic Centre (GC) region resulted from a series of lunar occultations in 1968 (Sandqvist 1971, 1973, 1974). They revealed that Sgr A consists of at least two components, now known as Sgr A East and Sgr A West (e.g. Downes & Martin 1971; Ekers et al. 1983), and that the region could be simply modeled by a rotating and contracting cloud of dust and molecules surrounding the continuum source.

Since then, interferometers have been playing an increasingly important role in the observations towards the GC. For example, searches for 1612-MHz OH maser emission associated with OH/IR stars (e.g. Lindqvist et al. 1992; Sjouwerman et al. 1998), and detections of individual 1720-MHz OH masers (Yusef-Zadeh et al. 1996) would not have been possible without interferometers because of confusion (Habing et al. 1983). Furthermore, in order to unravel the complex structure of Sgr A at higher resolution, Whiteoak et al. (1974) and Bieging (1976) used the Owens Valley interferometer to map the 6-cm H_2CO and 1667-MHz OH absorption, respectively. Subsequently, the Very Large Array (VLA) has been used in several configurations, with resolutions ranging from about $4''$ to an arcminute, to study the H I absorption (Liszt et al. 1983; Lasenby et al. 1989; Yusef-Zadeh et al. 1993; Zhao et al. 1995) and the H_2CO absorption (Whiteoak et al. 1983; Pauls et al. 1996). Magnetic field measurements have been performed using the Zeeman splitting of absorption lines in OH (Killeen et al. 1992) and H I (Plante et al. 1995).

The molecular complex associated with Sgr A has also been mapped in emission lines of many different molecules. Good reviews on the GC and its environment can be found in Morris & Serabyn (1996) and Mezger et al. (1996). The complex

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⋆ Figures 1 to 4 are only available in electronic form at http://www.edpsciences.org

Table 1. VLA 18-cm OH observational summary.

Frequency (MHz)	HPBW ($''\times''$)	PA ($^\circ$)	Resolution (km s^{-1})	$t_{\rm integration}$ (min)
1612	4.0×2.7	55.3	9.1	144
1665	3.9×2.9	64.7	8.8	169
1667	4.0×2.8	61.1	8.8	173
1720	3.7×2.7	57.1	8.5	148

consists predominantly of a molecular belt with density concentrations that define the "+50 km s^{-1} cloud" (M−0.02−0.07) and the "+20 km s^{-1} cloud" (M−0.13−0.08), as well as the Circumnuclear Disk (CND) (see e.g. Sandqvist 1989; Genzel 1989; Zylka et al. 1990; Lindqvist et al. 1995; Coil & Ho 2000). Other features seen towards the GC are the High Negative Velocity Gas (HNVG) (Yusef-Zadeh et al. 1993) and the Expanding Molecular Ring (EMR) (e.g. Güsten 1989). However, in contrast to emission of the total distribution of molecular gas, absorption only yields the foreground distribution of the molecule with respect to the continuum emission. Some indication of the relative location of the molecular region and the continuum sources can thus be obtained by comparing absorption and emission line data.

In this paper, we present arc-second resolution VLA observations of all four 18-cm OH lines with the BnA configuration. As the line intensity ratios of the four OH lines directly probe the population distribution of the gas, observations of all four lines provide a measure of the physical state of the gas, in particular deviations from thermal equilibrium. We will highlight the most striking results here in Sect. 3, although a detailed physical analysis of the data will be published in a forthcoming paper. Some of the preliminary results have been presented by Sandqvist et al. (1987, 1989).

2. Observations and data reduction

The BnA configuration observations were performed in late June 1986 and a summary of these observations is presented in Table 1. The observations were performed with 18 VLA antennas at RHC polarisation using a total bandwidth of 3.125 MHz, corresponding to a total velocity coverage of about 550 km s^{-1}, divided into 63 spectral channels. The centre line-of-sight velocity with respect to the Local Standard of Rest (LSR) was −38 km s^{-1} at each of the observed transitions. All references to velocities in this paper refer to line-of-sight velocities with respect to the LSR. We used the NRAO Astronomical Image Processing System (AIPS) for the data reduction. Extensive flagging of corrupted data was needed for the 1612-MHz data set due to interference of the Russian GLONASS satellite system. Standard calibration of the visibility phases was done against the nearby source B1748−253; standard amplitude and bandpass calibration were obtained using 3C 286.

Because the high-velocity tails of the OH main line transitions (1665 and 1667 MHz) overlap in the intermediate frequency range (i.e. around 1666 MHz), a direct subtraction of the continuum in the separate visibility data sets was not possible. Also, even when concatenating together the 1665- and 1667-MHz visibility data sets into one continuous frequency band, UVLIN could not be used due to different (u, v)-coverage. The continuum was therefore subtracted from the total image cube, made up from the unsubtracted 1665- and 1667-MHz image cubes concatenated together. The subtraction was done using IMLIN with a linear fit over feature-free channels on either side of the total 1665- and 1667-MHz absorption complex. For self-calibration of the phases we used the continuum image derived with IMLIN, partly because there were no strong masers, and partly because the continuum could not be subtracted from the visibility data. We made maps with 1024×1024 pixels of $0\farcs75$ with both the CLEAN as well as the maximum entropy algorithm (VTESS). Although the two methods gave similar results, the 103-channel image cube was faster to generate with the maximum entropy method (MEM) and of somewhat higher quality, i.e. less mottled. Typical rms noise levels are 5 mJy/beam. The resulting 96 MEM absorption channel maps can be found in Figs. 1 and 2, where some concatenation effects can be seen as noise, e.g. at the left side of the 1667-MHz maps at velocities of +100 to +140 km s^{-1} in Fig. 2. Note that although Figs. 1 and 2 are labeled with velocities for 1665 and 1667 MHz, respectively, Figs. 1 and 2 do form one continuous sequence.

In order to produce the maps of the absorption and line emission in the OH satellite line transitions (1612 and 1720 MHz), we subtracted out a linear fit over frequency to the continuum directly from the visibilities with UVLIN. A strong maser spot was then used to self-calibrate both the amplitudes and phases of the visibilities. For the 1612- as well as the 1720-MHz data we made 63 cleaned channel maps of 1024×1024 pixels of $0\farcs75$ with IMAGR. The relevant 48 channels roughly covering $+200 \gtrsim v_{\rm LSR} \gtrsim -200$ km s^{-1} are shown in Figs. 3 and 4. We also checked our CLEAN maps for 1612 and 1720 MHz by using the MEM algorithm for a small number of selected channels. No significant differences were found.

3. Results

Below we first discuss results on the OH absorption measurements; a second section discusses the OH emission features.

3.1. OH absorption results

The OH results, presented here in Figs. 1–4, have not been corrected for the effect of the background continuum distribution on the absorption distribution. For the simple case, where all the continuum is behind the absorbing clouds, the distribution of the line-to-continuum is the best indicator of the variation of the OH column density. This apparent opacity is given by $-T_{\rm L}/T_{\rm C}$ where $T_{\rm L}$ is the absorption line temperature which is a negative number, and $T_{\rm C}$ is the continuum temperature. Apparent opacity maps have been produced for several features and are presented below. The situation in the Sgr A Complex is, however, considerably more complicated since some of the continuum may be in front of some of the absorbing clouds. Furthermore, there is an additional problem inherent in interferometer observations: if the contiuum distribution is more

extended than the line absorption feature (which is true for the case of the Sgr A Complex), then the continuum will sometimes be more resolved out leading to incorrect continuum levels. This, in turn, leads to the possibility of $-T_L/T_C$ being greater than one, in some regions. This problem will require a certain amount of modelling and will be studied in more detail in a subsequent paper.

3.1.1. Molecular belt +50 and +20 km s^{-1}-clouds

The Sgr A +50 and +20 km s^{-1}-clouds appear to be massive condensations immersed in an extended molecular belt which stretches in a northeastern direction from south of the continuum Sgr A Complex to east of the Complex. While the cores of the two molecular clouds lie just outside the continuum Complex, the outer lower levels of the belt do coincide with the directions to the continuum components. The concept of a continuous belt is further strengthened by the gradual, although not linear, velocity gradient from positive velocities in the northeast through the +50 km s^{-1} and the +20 km s^{-1} clouds to negative velocities in the southwest (see e.g. Sandqvist 1989). The line-of-sight location of the different continuum and molecular components have been presented schematically by Coil & Ho (2000) and earlier by Sandqvist (1989) and Zylka et al. (1990).

It is apparent in all four lines, but especially in the high-resolution 1665- and 1667-MHz maps in Figs. 1 and 2, that the gas in the the molecular belt is seen clearly in absorption against the shell structure of Sgr A East – the eastern and most of the western parts of the shell – but not against the spiral structure of Sgr A West. See, for example, the sequence of velocity maps from +67.7 to +23.6 km s^{-1} in the 1667-MHz line in Fig. 2 (or in colour in Fig. 5; see also the $-T_L/T_C$ maps in Fig. 6) where the absorption can be followed across the major part of the Sgr A East shell. The dominant absorption moves westward as the velocity decreases, in agreement with the well-known velocity gradient which exists along the molecular belt. However, there is an obvious lack of absorption at the velocities around +32.4 and +23.6 km s^{-1} across Sgr A West and the expected western part of the Sgr A East shell at this position. Contrast this sequence with the +6.0 km s^{-1} map where at least the North Arm of Sgr A West is seen clearly being absorbed. The absorption in this latter map originates in a widespread overlying cloud well in front of Sgr A and the associated absorption distribution reflects the case where absorption is occurring against all the continuum in each direction. This may imply that a part of the molecular belt lies between the two continuum components, behind Sgr A West and in front of Sgr A East, implying in turn that Sgr A East is behind Sgr A West. Pedlar et al. (1989) and later investigations have presented strong evidence that Sgr A East lies behind Sgr A West. On the other hand, one would in this case still expect to see some absorption towards the Sgr A East western shell in the region of Sgr A West, and none is seen. There may therefore be a real distinct lack of absorbing OH gas in this region at the expected velocities. This missing gas may have been dispersed by the high velocities if Sgr A East and West are intertwined with the molecular belt. Alternatively, it may have been dissociated by the strong UV radiation field present in this region.

There seems to be very little absorption in the molecular belt towards Sgr A*, shown in the images of Figs. 1–4 as a cross, or in Figs. 5 and 6 as a tiny contour, at position (B1950.0) 17:42:29.33 −28:59:18.6, but here the picture is complicated by the OH-Streamer which is discussed in the next Sect. (3.1.2). Not until the velocity drops to −2.8 km s^{-1} is there clear unresolved absorption towards Sgr A*, and this can then be followed only through the −11.6 and −20.4 km s^{-1} maps which are not related to the molecular belt with its velocities of between +20 and +30 km s^{-1} at the position of Sgr A*. We believe that the absorption towards Sgr A* is really this localized and is related in some intriguing manner to the OH-Streamer. By −29.2 km s^{-1} the absorption has disappeared and does not appear again until the velocity region covered by the EMR, namely −126.0 to −143.6 km s^{-1}.

It is significant that no absorption at the molecular belt velocities is seen towards the compact H II regions lying just east of the Sgr A East shell, implying that these star formation regions must lie on the near (sunward) side of the molecular belt. Only in the 1667-MHz OH line at velocities near −126.0 km s^{-1} is there weak absorption towards these H II regions. This again is the velocity range of the EMR, so, at the very least, they are inside this molecular feature of the GC region. Their positions coincide quite well with the peak of the +50 km s^{-1} cloud (e.g. Sandqvist 1989), and Goss et al. (1985) have in fact detected H76α recombination lines from the H II regions with velocities in the range +43 to +52 km s^{-1}. These velocities are somewhat blueshifted with respect to the +50 km s^{-1} cloud, which would be in agreement with a model indicating their outflowing from the cloud, thus placing them on the near side of the cloud. This is also in agreement with the conclusions drawn by Cotera et al. (2000) from observations of interstellar extinction in the near infrared. Our OH observations would thus seem to support the association of the H II regions with the near side of the +50 km s^{-1} cloud.

A perusal of the satellite line maps in Figs. 3 and 4 shows that the maximum OH absorption towards Sgr A East is considerably deeper in the 1612-MHz line than in the 1720-MHz line, about 60 mJy/beam versus 45 mJy/beam, respectively. This phenomenon is generally true for most of the absorption seen in these maps.

3.1.2. The CND and the OH-Streamer

The Northeast and Southwest OH components of the CND seen already in the 1968 occultation maps (Sandqvist 1974) are very clear in the VLA BnA maps. They can be traced out to velocities of +150 and −150 km s^{-1}, respectively, in both the 1665- and 1667-MHz lines in Figs. 1 and 2. The signature of the rotating torus is especially clear southwest of Sgr A* in maps at velocities ranging from −20 to −100 km s^{-1} where the peak absorption intensities reach 89 and 106 mJy/beam for the 1665- and 1667-MHz OH lines, respectively. The northeastern part of the CND can be seen, for example, in the +94.1 to +120.5 km s^{-1} maps where the peak absorption intensities

are 68 mJy/beam for both the 1665- and 1667-MHz lines. While considerably weaker, the CND can still be traced in the 1612-MHz OH line in Fig. 3, and even in the 1720-MHz OH line in Fig. 4 although only then with some difficulty.

A unique feature – which we called the "1667-MHz OH Streamer" in our preliminary papers (Sandqvist et al. 1987, and Sandqvist et al. 1989) – is now seen to have been detected in three of the four OH lines, namely the 1612-, 1665- and 1667-MHz lines. It is not detectable in the 1720-MHz line. This molecular gas streamer stretches from the CND's southwestern region inwards through the cavity, inside the CND (at least in projection), to the compact non-thermal radio source, Sgr A*.

The OH-Streamer can best be followed in the 1667-MHz maps of Fig. 2. We also present the six most important channel maps (+67.6 to +23.6 km s^{-1}) of Fig. 2 in colour in Fig. 5. It may first be (marginally) distinguished from other features at a velocity of +76.4 km s^{-1}. At a velocity of +67.6 km s^{-1} it is seen to cut across the southwestern part of the CND and into the cavity. It stands out clearest at a velocity of +58.8 km s^{-1}. Note that the negative velocity of the CND in this region differs from that of the OH-Streamer by more than 100 km s^{-1}, and the OH-Streamer may not necessarily be in the same plane as the CND. At these velocities, the "head" of the OH-Streamer (the part closest to Sgr A*) does *not* coincide with Sgr A* but lies slightly to its northwest. The OH-Streamer then "shrinks" and moves closer to Sgr A* as the velocity drops until it coincides with Sgr A* near +23.6 km s^{-1}. Its coincidence with Sgr A* can then be followed through to the −20.4 km s^{-1} map, after which there is no more detectable absorption towards Sgr A*.

It is intriguing to follow the structure of the "head" of the OH-Streamer as its trailing and shrinking part sweeps in a counterclockwise direction while the velocities drop from +50.0 to −20.4 km s^{-1}. From pointing towards the "tail" of the OH-Streamer in a southwest direction at +50 km s^{-1}, the trailing end of the head points west at +41.2 km s^{-1}, north at +32.4 and +23.6 km s^{-1}. This gradual movement inwards and final coincidence with Sgr A* makes us believe that the OH-Streamer does in fact (and not just in projection) exist in the cavity inside the CND and interacts with Sgr A*. An orbital model for the OH-Streamer will be developed and presented in a future paper.

We have produced $-T_L/T_C$ maps for the 1667-MHz OH line in the velocity range +67.7 to +23.6 km s^{-1} and present them in colour in Fig. 6. Here we have assumed that all the contiuum is behind the OH gas. The effect of the velocity gradient across the molecular belt, discussed in Sect. 3.1.1, is seen also in the $-T_L/T_C$ maps. The OH-Streamer retains its shape and general behaviour although it may be somewhat wider than was apparent in the line intensity maps in Fig. 5. The assumption that the OH-Streamer is in front of the southwest arm of Sgr A West may thus be correct. Its $-T_L/T_C$ value is considerably lower, though, than the gas in the molecular belt seen against the shell of Sgr A East.

The OH-Streamer can also be traced in the same velocity maps of the 1665-MHz OH line in Fig. 1. In the 1612-MHz OH line maps of Fig. 3, the OH-Streamer is detectable at +61.9 km s^{-1} and also seen at +71.0 and +52.8 km s^{-1}. We were not able to recognize any absorption by the OH-Streamer in the 1720-MHz line (Fig. 4). The relative intensities of the absorptions in the two main lines are not far from those expected (5:9) from an optically thin gas in LTE. In the head of the OH-Streamer, near Sgr A*, the ratio of the 1665:1667-MHz line intensities, 42 mJy/beam:79 mJy/beam, is about 5:9.4. The ratio increases slightly to about 5:7.5 (20 mJy/beam:30 mJy/beam) in the tail which extends southwestward towards the CND, possibly indicating the onset of 1667-MHz line saturation caused by higher density in this region. The $-T_L/T_C$ maps in Fig. 6 also indicate higher densities in the tail than in the head. However, the absorption intensities of the 1612-MHz line are 20 and 12 mJy/beam in the head and tail, respectively, which is surprisingly high. The excitation of the lines is obviously complex and requires deeper study. It is also difficult to understand how the OH-Streamer can survive in the strong UV radiation field in the cavity inside the CND, but it may be very young and the OH may be protected to some extent if it is intermingled with dust.

3.1.3. Other absorption features in the BnA maps

In Sect. 3.1.1, we briefly mentioned the absorption towards Sgr A* at velocities near −126.0 to −143.6 km s^{-1} caused by the EMR. At a velocity of −152.4 km s^{-1} there is practically no OH absorption, but starting at −161.3 km s^{-1}, reaching a maximum near −178.9 km s^{-1} and petering out near −222.9 km s^{-1} there is extended OH absorption in the region of Sgr A West. This HNVG structure is apparently double-lobed as seen in the OH gas, one lobe to the northeast of Sgr A*, the other to the southwest of Sgr A*, somewhat similar to the orientation of the inner part of the CND. This double-lobed appearance is retained even in $-T_L/T_C$ maps, an example of which is shown in Fig. 7. Only the southwest lobe, though, is seen at the very negative velocities between −196.5 and −222.9 km s^{-1} (see Fig. 2). When studying a feature in the overlapping velocity range of the 1665- and 1667-MHz lines, confusion can be mostly eliminated by studying the corresponding feature in the high negative velocity range of the 1667-MHz line and the high positive velocity range of the 1665-MHz line, respectively. The southwest HNVG can also be seen in the −165.6 and −174.1 km s^{-1} maps of the 1720-MHz OH line and the northeast part weakly in the 1612-MHz line at −174.2 km s^{-1}. The HNVG, as seen in H I by Yusef-Zadeh et al. (1993) is predominantly distributed along the rotation axis of the CND. A detailed model for the HNVG has been developed by Zhao et al. (1995).

3.2. OH emission results

The four sets of "maxmaps" and "minmaps" of the data cubes are shown in Figs. 8 and 9. The maxmap/minmap presentation method was developed by Sjouwerman (1997). The maxmaps (or minmaps) are produced through projection of the three-dimensional data cube onto a two-dimensional sky image map. This is done by storing the maximum (or minimum) intensity over the whole velocity axis for each pixel in the sky

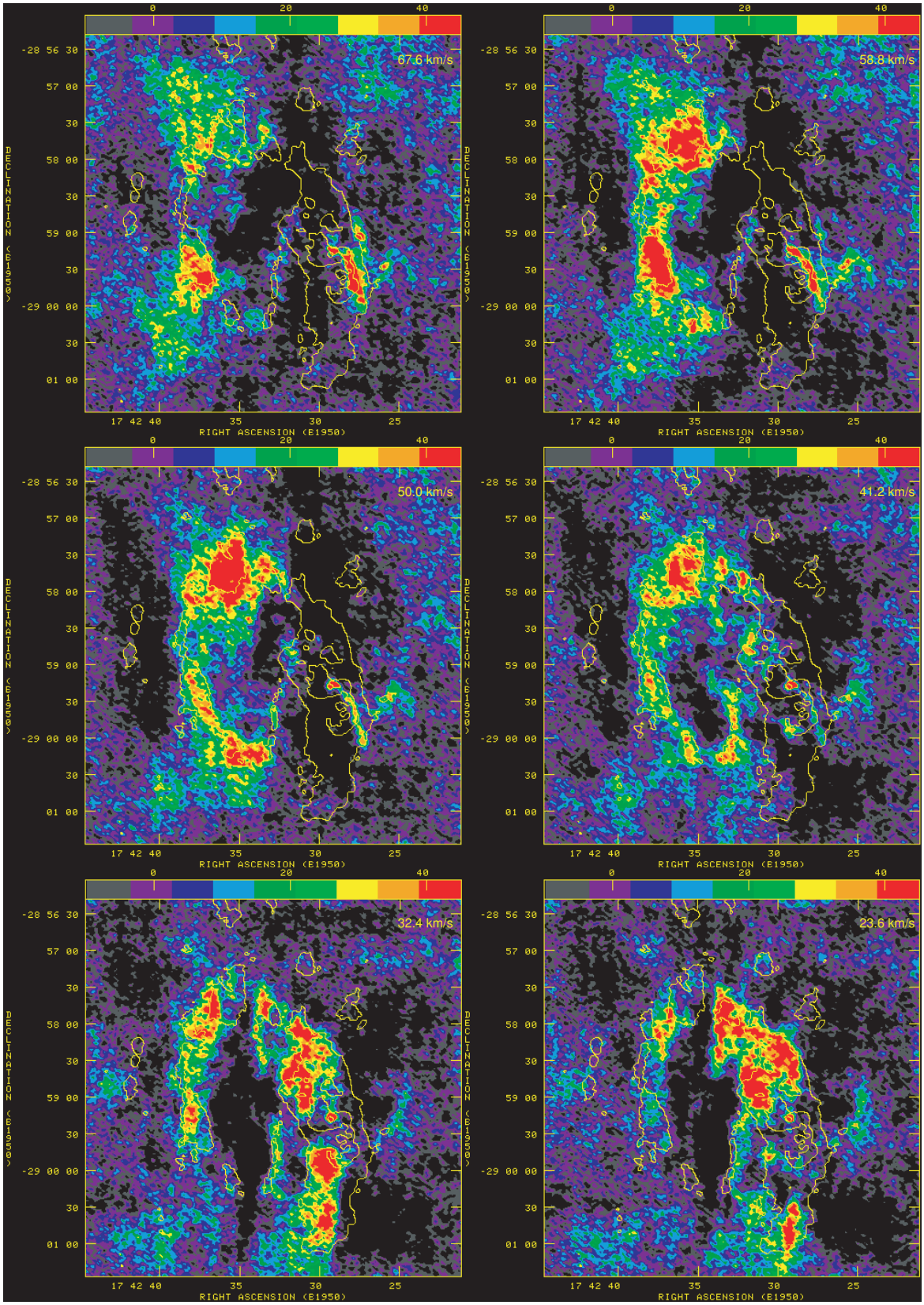

Fig. 5. 1667-MHz OH absorption, highligting in colour the OH-Streamer near Sgr A* and the absorption towards Sgr A East. The range in colour is in mJy/beam. The velocities are indicated in the upper right corner. The contour level of 15 mJy/beam indicates Sgr A East and the three compact H II-regions to its east, the contour level 93 mJy/beam indicates Sgr A West, and the contour level 500 mJy/beam indicates Sgr A*.

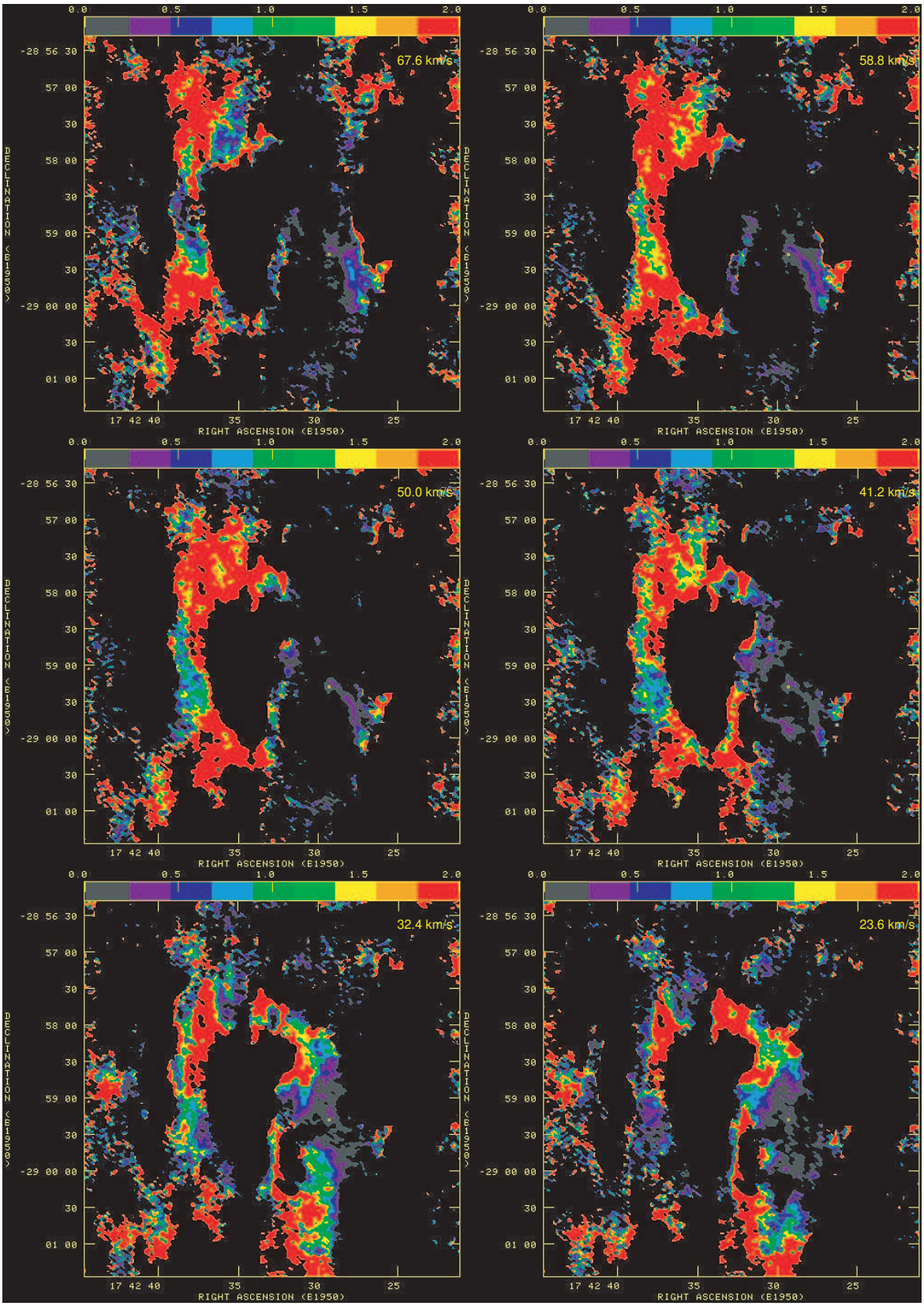

Fig. 6. 1667-MHz OH line-to-continuum distribution ($-T_L/T_C$), highligting in colour the OH-Streamer near Sgr A* and the absorption towards Sgr A East. The range of the colour scale is from 0.0 to 2.0. The velocities are indicated in the upper right corner. The position of Sgr A* is indicated by a tiny yellow contour.

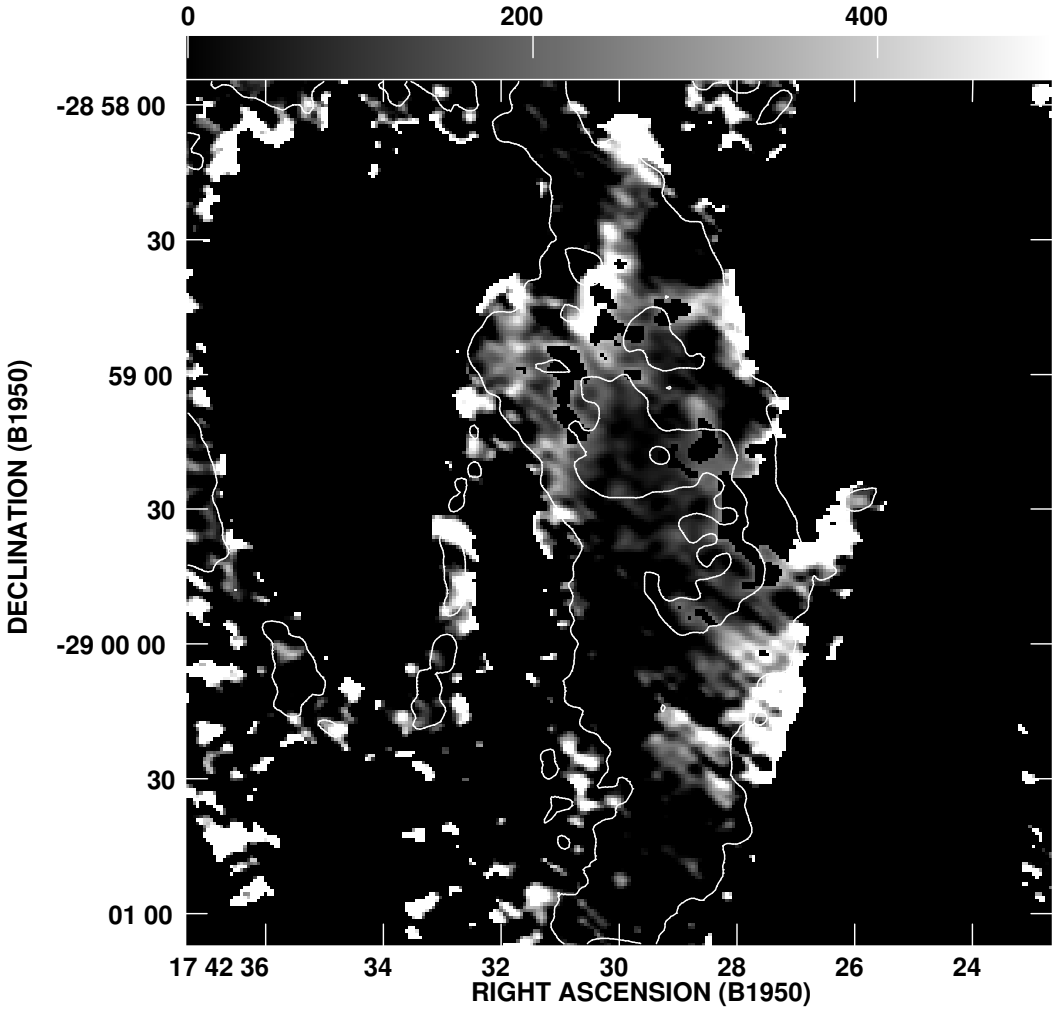

Fig. 7. 1667-MHz OH line-to-continuum distribution ($-T_L/T_C$) for the High Negative Velocity Gas (HNVG) at $v_{LSR} = -170.1$ km s^{-1}. The range of the gray scale is from 0 to 0.500. The contour level of 15 mJy/beam indicates the western shell of Sgr A East, the contour level 93 mJy/beam indicates Sgr A West, and the contour level 500 mJy/beam indicates Sgr A*.

plane. Maxmaps and minmaps can nowadays easily be generated with the AIPS task SQASH. The maxmaps are useful tools for identifying masers and possible regions of extended emission, while the minmaps give an overall view of the spatial distribution of the absorption.

Weak extended emission can be seen in the main line maxmaps in Fig. 8. Although weak, extended and diffuse OH emission has been reported elsewhere in the Galactic Centre region, in the 1720-MHz satellite line (Yusef-Zadeh et al. 1999), and in other galaxies in the 1667-MHz line (e.g. Pihlström et al. 2001), the extended OH line emission features in our observations may be indicative of an instrumental effect. We consider it unlikely to have introduced this weak emission by using a non-standard continuum subtraction for the main line OH data. We do not see any *extended* emission regions in the satellite lines (Fig. 9), where we have also carefully checked especially individual 1720-MHz maps with velocities close to those of narrow emission features in single-dish integrated spectra (e.g. Whiteoak & Gardner 1976).

It is known that lack of zero- and short-spacing observations in interferometric data may cause spurious emission or absorption features in interferometer maps. The lack of visibility data from short baselines results in an unsampled central hole in the (u, v)-coverage, which cannot be reconstructed correctly by the deconvolution algorithms without a-priori information (such as single-dish observations). In our case, this results in an insensitivity to structures larger than about 80″. This effect may also lead to a (negative) bowl-like structure around strong emission in the central parts of the brightness distribution. The velocity line maps and the background continuum are both affected, although somewhat offset in position from each other depending upon their relative intensity distribution. Ghost line emission features may therefore occur after continuum subtraction. The negative bowl appears to be of the order of 10 mJy/beam, i.e. about 15% of the extended emission of the Sgr A Complex.

3.2.1. Point-like OH masers

A large number of maser emission point sources are visible in the satellite lines of 1612- and 1720-MHz OH maxmaps in Fig. 9. Some 1665- and 1667-MHz OH main line maser point

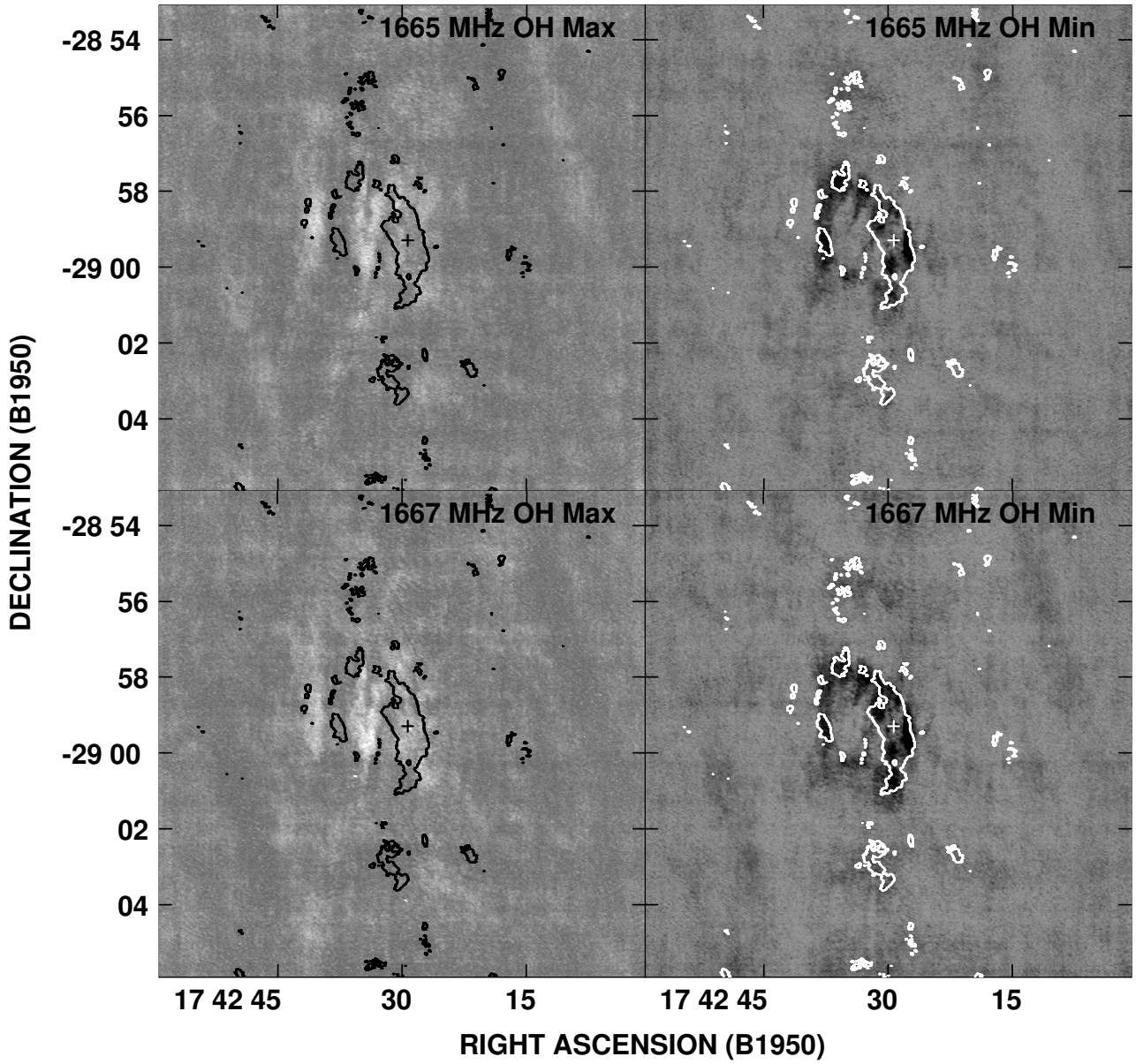

Fig. 8. 18-cm OH main line summary; Maxmaps and Minmaps at 1665 MHz and 1667 MHz. Continuum contours and a cross at the position of Sgr A* are placed to guide the eye. The gray scale showing the relative noise of the maps is set to match that of the satellite lines (see text).

sources are also visible in Fig. 8. The entire list of emission point source detections is presented in Table 2. Although these data were taken in 1986, the 1612-MHz maser sources have already been published by Lindqvist et al. (1992; L-92) and subsequently by Sjouwerman et al. (1998; S-98), and most of the 1720-MHz masers by Yusef-Zadeh et al. (1996; YZ96).

There are four new 1720-MHz OH masers in Table 2, of which the most interesting may be OH359.930−0.049 and which meanwhile also has been noticed in this 1986 data set, and commented on by Yusef-Zadeh et al. (2001; YZ01). Its $v_{\rm system}$ of -132 km s^{-1} and its position in the southwestern part of the CND makes it an important symmetric counterpart to the YZ96 maser "B" with its $v_{\rm system}$ of $+132$ km s^{-1} in the northeastern part of the CND. It suggests a mass enclosed within about 47.″4 (1.84 pc) of Sgr A*of at least 7.5×10^6 $M_\odot$, assuming Keplerian motion in an edge-on disk at 8 kpc distance. Granting up to 3.7×10^6 $M_\odot$ for the black hole at the position of Sgr A* (Schödel et al. 2002), this would then imply that at least 50% of the total mass within the CND is contained in the enclosed stellar cluster and molecular cloud complexes. As expected, none of the 1720-MHz (shock-excited) masers coincide with the circumstellar 1612-MHz (FIR pumped) masers.

Four new 1667-MHz OH masers are presented in Table 2, all of them identified as circumstellar masers found by Lindqvist et al. (1992) in the 1612-MHz line. Two of them, corresponding to the strongest 1612-MHz masers, are also newly detected in the 1665-MHz line. Although we are dealing with (mostly) intrinsic variable sources, we can compare the data from the different transitions for each star because the data sets are taken simultaneously. Several interesting observations

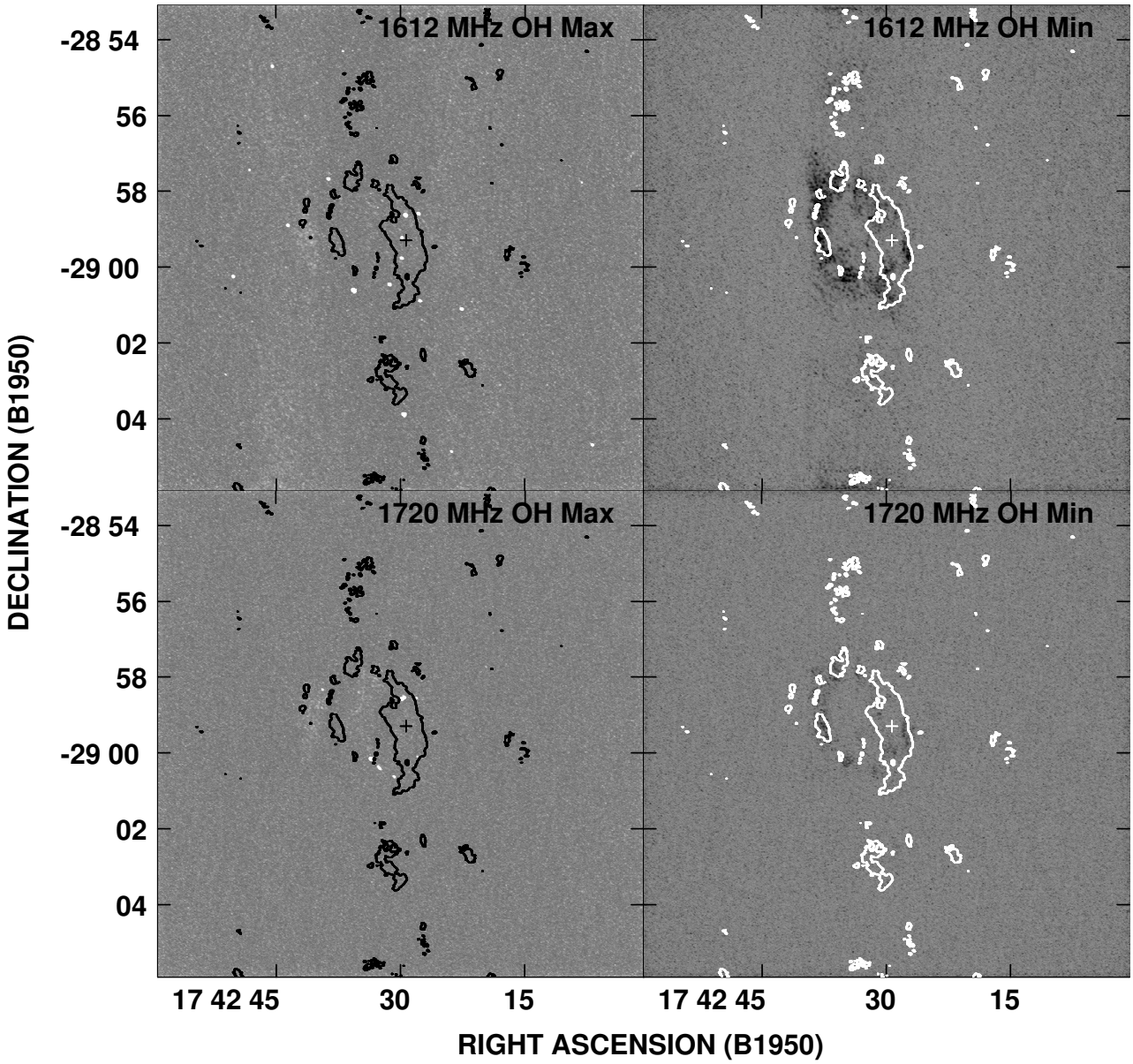

Fig. 9. 18-cm OH satellite line summary; Maxmaps and Minmaps at 1612 MHz and 1720 MHz. Continuum contours and a cross at the position of Sgr A* are placed to guide the eye. The gray scale showing the relative noise of the maps is set to match that of the main lines (see text).

can already be made, although one should be cautious with the small numbers involved.

First, the 1665-MHz masers are only found in the strongest 1612-MHz masers at a level of 1.5 and 2.8% with respect to the 1612-MHz flux. If this ~2% is a characteristic fraction, and recalling that the noise in the maps is about 5 mJy/beam, then indeed these two sources are the only ones that would exceed five times the noise level at 1665-MHz.

Second, the two sources that show 1665-MHz masers (i.e. the strongest 1612-MHz sources OH359.880−0.087 and OH359.938−0.078) also show 1667-MHz maser emission at a level between 4.4 and 9.5% of the 1612-MHz flux. Inspecting the rotational transition levels of OH, it becomes clear that in OH359.952−0.036 and OH359.977−0.087, the 1612-MHz and 1667-MHz masers are competing to populate the $J = 3/2$, negative parity $F = 2$ level. That OH359.880−0.087 and OH359.938−0.078 lack in relative 1667:1612-MHz intensity might be due to the influence of the 1665-MHz line, where the relative 1667:1612-MHz intensity in particular seems to decrease with the stronger 1665-MHz maser.

As a final remark, although OH359.952−0.036 and OH359.977−0.087 are "normal" 1612-MHz maser emitters, none of the other "normal" 1612-MHz masers that would have about 13% of the 1612-MHz intensity above the noise (25 mJy/beam) has been found in the 1667-MHz line. Except for a very weak (18.67 mag) K-band 2.2 μm counterpart for OH359.880−0.087 measured by Blommaert et al. (1998), none of the 1667-MHz masing stars have been found in the K-band by either Blommaert et al. (1998), nor by Glass et al. (2001). In that respect the 1667-MHz masers are not normal – they are "redder" than the others. It should also be noted that OH359.938−0.078 and OH359.952−0.036 are not variable in

Table 2. Point source detections.

Name	Position in B1950 (RA)	(Dec)	Position in J2000 (RA)	(Dec)	$v_{\rm system}$ (km s^{-1})	Flux (mJy)	Reference
OH 1612 MHz masers							
OH359.825−0.024	17 42 06.961	−29 04 40.90	17 45 17.824	−29 05 52.10	−56	190	L-92 (47)
OH359.924+0.034	17 42 08.300	−28 57 44.95	17 45 18.987	−28 58 56.05	−83	72	L-92 (60)
OH359.990+0.030	17 42 18.210	−28 54 32.29	17 45 28.817	−28 55 42.66	−65	64	S-98
OH359.906−0.036	17 42 21.493	−29 00 54.62	17 45 32.262	−29 02 04.73	−47	40	S-98
OH359.888−0.051	17 42 22.486	−29 02 19.15	17 45 33.292	−29 03 29.21	−56	62	L-92 (53)
OH359.906−0.041	17 42 22.714	−29 01 05.75	17 45 33.488	−29 02 15.76	−138	284	L-92 (58)
OH359.855−0.078	17 42 23.935	−29 04 51.52	17 45 34.805	−29 06 01.43	+3	129	L-92 (50)
OH359.911−0.059	17 42 27.482	−29 01 22.96	17 45 38.264	−29 02 32.61	−79	54	L-92 (59)
OH359.918−0.055	17 42 27.700	−29 00 53.27	17 45 38.469	−29 02 02.89	−292	299	S-98
OH359.952−0.036	17 42 27.897	−28 58 35.23	17 45 38.608	−28 59 44.83	+85	289	L-92 (66)
OH359.954−0.041	17 42 29.463	−28 58 37.61	17 45 40.175	−28 59 47.09	+71	731	L-92 (67)
OH359.880−0.087	17 42 29.641	−29 03 52.92	17 45 40.486	−29 05 02.37	−20	1085	L-92 (52)
OH359.946−0.047	17 42 29.892	−28 59 13.47	17 45 40.619	−29 00 22.89	−24	145	L-92 (65)
OH359.939−0.052	17 42 29.919	−28 59 45.56	17 45 40.659	−29 00 54.97	+53	256	L-92 (63)
OH359.932−0.059	17 42 30.575	−29 00 16.96	17 45 41.329	−29 01 26.38	−156	44	S-98
OH359.932−0.063	17 42 31.383	−29 00 27.57	17 45 42.141	−29 01 36.86	−92	286	L-92 (61)
OH359.970−0.049	17 42 33.681	−28 58 03.05	17 45 44.381	−28 59 12.31	+89	39	L-92 (69)
OH359.985−0.042	17 42 34.034	−28 57 04.34	17 45 44.707	−28 58 13.46	−11	28	S-98
OH359.938−0.077	17 42 35.685	−29 00 36.57	17 45 46.447	−29 01 45.54	−83	2408	L-92 (62)
OH359.902−0.103	17 42 36.608	−29 03 13.52	17 45 47.439	−29 04 22.56	−102	44	S-98
OH 0.030−0.026	17 42 36.999	−28 54 17.89	17 45 47.602	−28 55 26.76	−56	70	S-98
OH359.986−0.061	17 42 38.742	−28 57 39.90	17 45 49.430	−28 58 48.63	+16	246	L-92 (73)
OH359.980−0.077	17 42 41.720	−28 58 27.08	17 45 52.428	−28 59 35.59	+89	40	S-98
OH359.977−0.087	17 42 43.616	−28 58 53.52	17 45 54.336	−29 00 01.88	+12	220	L-92 (72)
OH 0.040−0.056	17 42 45.477	−28 54 42.46	17 45 56.091	−28 55 50.68	+71	147	L-92 (81)
OH 0.007−0.088	17 42 48.232	−28 57 26.00	17 45 58.915	−28 58 34.00	+3	76	L-92 (77)
OH 0.053−0.062	17 42 48.853	−28 54 14.07	17 45 59.458	−28 55 22.21	+7	26	L-92 (82)
OH359.971−0.119	17 42 50.180	−29 00 14.39	17 46 00.935	−29 01 22.24	−11	227	L-92 (70)
OH359.936−0.144	17 42 51.127	−29 02 49.22	17 46 01.947	−29 03 56.99	−11	61	L-92 (7-s)
OH 1665 MHz masers							
OH359.880−0.087	17 42 29.628	−29 03 52.91	17 45 40.472	−29 05 02.27	−29	30	
OH359.938−0.077	17 42 35.678	−29 00 36.25	17 45 46.440	−29 01 45.16	−91	35	
OH 1667 MHz masers							
OH359.952−0.036	17 42 27.907	−28 58 35.21	17 45 38.618	−28 59 44.75	+63	37	
OH359.880−0.087	17 42 29.597	−29 03 53.40	17 45 40.442	−29 05 02.80	−29	103	
OH359.938−0.077	17 42 35.682	−29 00 36.55	17 45 46.444	−29 01 45.50	−82	106	
OH359.977−0.087	17 42 43.647	−28 58 53.57	17 45 54.367	−29 00 01.94	−3	29	
OH 1720 MHz masers							
OH359.930−0.049	17 42 28.017	−29 00 05.83	17 45 38.764	−29 01 15.27	−132	70	YZ01
OH359.952−0.036	17 42 28.181	−28 58 31.47	17 45 38.888	−28 59 40.89	+47	118	YZ96 (C)
OH359.955−0.040	17 42 29.666	−28 58 32.18	17 45 40.374	−28 59 41.48	+132	290	YZ96 (B)
OH359.960−0.037	17 42 29.705	−28 58 09.87	17 45 40.403	−28 59 19.16	+30	37	
OH359.955−0.042	17 42 29.913	−28 58 34.44	17 45 40.621	−28 59 43.71	+132	658	YZ96 (B)
OH359.928−0.062	17 42 30.843	−29 00 37.40	17 45 41.603	−29 01 46.59	+56	131	YZ96 (G)
OH359.934−0.066	17 42 32.616	−29 00 23.40	17 45 43.371	−29 01 32.45	+56	248	YZ96 (D,E,F)
OH359.939−0.067	17 42 33.588	−29 00 09.07	17 45 44.337	−29 01 18.04	+64	1290	YZ96 (A)
OH359.966−0.056	17 42 34.878	−28 58 28.00	17 45 45.586	−28 59 37.03	+39	75	
OH359.977−0.069	17 42 39.462	−28 58 19.67	17 45 50.165	−28 59 28.20	+56	93	

their 1612-MHz maser emission, whereas OH359.977−0.087 has a, for the GC region, long period of 1070 days. The source with a K-band detection, OH359.880−0.087, has a more typical period of 760 days (the periods are taken from van Langevelde et al. 1993). It is tempting to conclude that these sources are evolving away from the OH/IR star stage to

become planetary nebulae, where the non-variable 1612-MHz sources OH359.938−0.078 and OH359.952−0.036 are the most transformed sources, and where OH359.880−0.087 has just started the transition.

4. Summary

We have performed observations of the four ground rotational state OH transitions towards the Galactic Centre with the VLA in the BnA configuration. Presented here are velocity maps of OH absorption at 9 $\mathrm{km\,s^{-1}}$ intervals and with angular resolution of $\approx 4''$. Maps of the $-T_\mathrm{L}/T_\mathrm{C}$ distribution for selected regions have confirmed and strengthened the reality of features and phenomena revealed in the absorption line maps. In a forthcoming paper, we shall determine relative intensities and optical depths of the structural features in all four ground-state lines, as well as perform a more detailed physical analysis of this large amount of data.

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